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# Sub-10 $\mu\text{m}$ grain size, $\text{Ba}_{1-x}\text{Ca}_x\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$ ( $x=0.10$ and $x=0.15$ ) piezoceramics processed using a reduced thermal treatment

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## Abstract

The solid-state synthesis of  $\text{Ba}_{1-x}\text{Ca}_x\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$  ( $x=0.10, 0.15$ ) (BCTZ) powder and the processing method of ceramics, by the use of reduced synthesis time and temperature (1250 °C for 2 h), are reported. Homogeneous and dense ( $\geq 95\%$ ) ceramic microstructures with sub-10  $\mu\text{m}$  grain size were obtained under all sintering conditions. A comparative study of their ferro-piezoelectric properties as a function of sintering temperatures is presented. The study shows the role of the grain size effect for improving both piezoelectric and ferroelectric properties of these materials. With an increase of the sintering temperature, grain growth was promoted; therefore, higher ferro-piezoelectric values were obtained (at 1400 °C, for  $x=0.10$ :  $d_{33}=300$  pC/N,  $d_{31}=-150$  pC/N,  $k_p=48\%$ ; for  $x=0.15$ :  $d_{33}=410$  pC/N,  $d_{31}=-154$  pC/N,  $k_p=50\%$ ). In addition, a diffuse phase transition is observed in these BCTZ ceramics with a Curie temperature near 100 °C at 1 kHz.

Keywords: BCTZ solid solutions, ferroelectrics, piezoelectrics

(Some figures may appear in colour only in the online journal)

## 1. Introduction

Nowadays, and due to the increasing importance of environmental protection, research on piezoelectric and ferroelectric materials has been focused on lead-free ceramics development, for which some objectives of industrial transference have already been achieved, whereas some basic problems remain still unsolved [1]. Barium titanate (BT)

binary systems [2–5] have been widely studied due to their non-toxicity and aim to overcome the main issues for practical applications, i.e., its low Curie temperature ( $T_C \approx 120$  °C) and electromechanical coupling factor ( $k_p \approx 35\%$ ).

BT undergoes three known structural phase transformations. Substituting the A and/or B sites of the BT perovskite structure makes it possible to modify its phase transition

temperature. The addition of  $\text{Ca}^{2+}$  into the  $\text{Ba}^{2+}$  site increases the stability temperature range of the tetragonal phase [2] by displacing the rhombohedral (R) to the orthorhombic (O) ( $T_1$ ) and the orthorhombic (O) to the tetragonal (T) ( $T_2$ ) ferroelectric-ferroelectric (FE) phases' transition temperatures. An optimum amount of  $\text{Ca}^{+2}$  ( $0.06 < x < 0.10$ ) enhances the dielectric properties of BCTZ materials [6]. The substitution of  $\text{Zr}^{4+}$  into  $\text{Ti}^{4+}$  sites in  $\text{Ba}(\text{Ti}_{1-y}\text{Zr}_y)\text{O}_3$  produces the shift toward a lower temperature of the tetragonal to cubic phase transition, i.e., ferroelectric to paraelectric transition at  $T_C$ , while the other two FE to FE phases' transition temperatures ( $T_1$  and  $T_2$ ) increase at different rates. For  $y=0.15$ , three ferroelectric phases (R, O and T) coexist near room temperature. It was expected to see a great polarizability at the pinched point; however, no ferroelectric and piezoelectric improvement was achieved. This was probably due to the degradation of ferroelectricity and weak ferroelectric-relaxor behavior at this level of  $\text{Zr}^{4+}$  content [3]. Furthermore, the ceramic shows broad dielectric peaks with frequency dispersion, i.e., ferroelectric relaxor behavior with increasing the  $\text{Zr}^{4+}$  concentration ( $y > 0.25$ ) [4, 5]. Recent work suggested a phase coexistence between  $0.02 > y > 0.05$  and  $0.07 < y < 0.09$  due to two polymorphic phase transitions, resulting in enhanced piezoelectric properties ( $225 \text{ pC/N} < d_{33} < 295 \text{ pC/N}$ ) [7].

Lately, pseudo-binary systems with high piezoelectric sensitivity  $(\text{Ba,Ca})(\text{Zr,Ti})\text{O}_3$  have been successfully developed as promising candidates for lead-free piezoelectric ceramics [8–10]. Among the drawbacks that exist in compositions with the highest piezoelectric activity ( $x=0.15$   $y=0.10$ ;  $d_{33} \approx 620 \text{ pC/N}$ ) are the low  $T_C \approx 93 \text{ }^\circ\text{C}$ , high synthesis ( $1300 \text{ }^\circ\text{C}$  for 2 h) and sintering ( $1450 \text{ }^\circ\text{C}$  for 3 h) conditions (not suitable for industrial production) and the subsequently large grain structures. Microstructures with average grain size  $> 10 \mu\text{m}$  obtained to date are not optimized for the present trend of device miniaturization. This feature result is important in order to increase the frequency range of the ultrasonic transduction. Therefore, strategies to reduce the synthesis and sintering temperatures are needed for processing routes towards sub- $10 \mu\text{m}$  grain size ceramics.

Some efforts have been dedicated to reducing the materials' processing temperature [11] using conventional ceramic methods and employing different modifying agents, i.e.,  $\text{Ba}_{0.85}\text{Ca}_{0.15}\text{Ti}_{0.90}\text{Zr}_{0.10}\text{O}_3$  ceramics with  $\text{CeO}_2$  added, calcined at  $1250 \text{ }^\circ\text{C}$  for 2 h and sintered at  $1350 \text{ }^\circ\text{C}$  for 4 h, showed a piezoelectric coefficient  $d_{33} \approx 600 \text{ pC/N}$  [12]. Moreover,  $\text{Ba}_{0.85}\text{Ca}_{0.15}\text{Ti}_{0.90}\text{Zr}_{0.10}\text{O}_3$  ceramics with  $\text{MnO}$  added, calcined at  $1100 \text{ }^\circ\text{C}$  for 4 h and sintered at  $1350 \text{ }^\circ\text{C}$  for 2 h, exhibited a piezoelectric coefficient  $d_{33} \approx 373 \text{ pC/N}$  [13]. Besides, LiF doping was used to reduce the sintering temperature at the expense of lower piezoelectric properties [14]. The authors recently achieved good piezoelectric properties in ceramics processed below  $1300 \text{ }^\circ\text{C}$  using the Pechini method [15].

In this work, we explore the possibility of optimizing the ceramic microstructure by a conventional solid-state technique using moderate sintering conditions, while keeping piezoelectric properties high enough for practical uses.

Investigations were carried out for  $\text{Ba}_{1-x}\text{Ca}_x\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$  ceramics with  $x=0.10$  and  $0.15$ , since their good piezoelectric properties were recently reported using conventional processing routes [8, 9, 16, 17]. The relationship between structure and electrical properties is discussed. The effects of the sintering conditions over structure and ferro-piezoelectric properties are also reported.

## 2. Experimental details

$\text{Ba}_{1-x}\text{Ca}_x\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$  ( $x=0.10$  and  $0.15$ ) ceramics were prepared by solid-state synthesis. The starting materials were  $\text{BaCO}_3$  (99.0%, Analytica),  $\text{CaCO}_3$  (99.0%, Fluka),  $\text{TiO}_2$  (99.0%, Aldrich) and  $\text{ZrO}_2$  (99.0%, Riedel-deHaën). Stoichiometric amounts of reagents were weighted and mixed with acetone for 30 min. They were calcined at  $1250 \text{ }^\circ\text{C}$  for 2 h. Afterwards, powders were ball milled for 24 h with zirconia ball media and alcohol, pressed into  $\sim 13.0 \text{ mm}$  diameter and  $\sim 1.5 \text{ mm}$  thickness pellets and sintered at  $1300 \text{ }^\circ\text{C}$ ,  $1350 \text{ }^\circ\text{C}$  and  $1400 \text{ }^\circ\text{C}$  for 2 h.

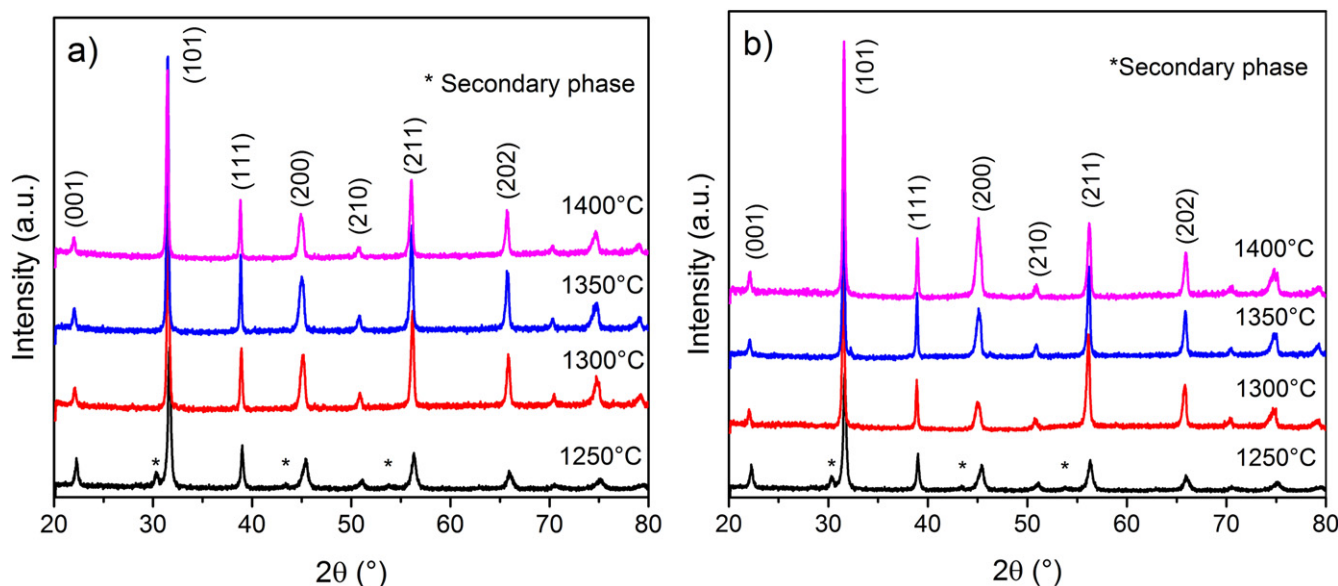
The phase structure of BCTZ ceramics was analyzed using x-ray diffraction with a  $\text{Cu } k\alpha 1$  radiation ( $1.54178 \text{ \AA}$ , Bruker D8 Advance with a  $0.016^\circ$  step size in  $2\theta$  and 1 s of integration time). Morphology of the samples was studied by a field emission scanning electron microscope (SEM) (JEOL J7600f) and computer-aided image analysis with Image J software. The average grain size of a large sample area was determined by the interception method.

Samples used for piezoelectric and dielectric characterization were polished and silver paste was applied on both sides of the samples at  $700 \text{ }^\circ\text{C}$  for 10 min to form electrodes.

Dielectric permittivity ( $\epsilon_{33}^T$ ) and losses ( $\tan \delta$ ) were measured with a precision impedance analyzer (Agilent 4294A) from room temperature up to  $180 \text{ }^\circ\text{C}$ . Then, ceramics were poled at  $2 \text{ kV mm}^{-1}$  for 30 min at room temperature. Piezoelectric constants  $d_{33}$  and  $-d_{31}$  were measured using a wide-range  $d_{33}$ -meter (APC International). Electromechanical coupling factor  $k_p$ , elastic compliances  $s_{11}^E$ ,  $s_{12}^E$  and piezoelectric constants  $-g_{31}$ ,  $g_{33}$  were calculated using the resonance method and an automatic iterative analysis method of complex impedance at the radial mode of the disks [18]. The piezoelectric properties were measured after 24 h of the poling process. Ferroelectric hysteresis loops were measured at room temperature using a Radiant workstation at  $100 \text{ Hz}$ .

## 3. Results and discussion

Figure 1 shows the x-ray diffraction (XRD) patterns of  $\text{Ba}_{1-x}\text{Ca}_x\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$  ( $x=0.10$  and  $0.15$ ) calcined powders at  $1250 \text{ }^\circ\text{C}$  for 2 h and those of sintered ceramics at different temperatures. The powders have traces of secondary phases, which indicates the fact that all reactions were not completed at this temperature and they are still reactive upon further thermal treatments, thus enhancing sintering. The samples showed a single perovskite phase crystal structure, and the diffraction peaks of the samples were



**Figure 1.** XRD patterns of BCTZ ceramics synthesized at 1250 °C and sintered at 1300, 1350, 1400 °C: (a)  $x=0.10$  and (b)  $x=0.15$ .

indexed to the perovskite-type cubic prototype structure. XRD patterns indicate that  $\text{Ca}^{2+}$  and  $\text{Zr}^{4+}$  ions incorporate to the  $\text{BaTiO}_3$  lattice to form a solid solution.

The crystal structure of pure  $\text{BaTi}_{0.9}\text{Zr}_{0.10}\text{O}_3$  at room temperature is rhombohedral [3], which is indicated by the single (200) peak, the doublet (111)/(-111) and the small (311) peak. With the increase of  $\text{Ca}^{2+}$  content, pronounced splitting of the (002) peak from the (200) peak has been observed, demonstrating a phase transition from the rhombohedral to tetragonal phase [16], which is stabilized at room temperature by the substitution of  $\text{Ba}^{2+}$  by  $\text{Ca}^{2+}$ . Then, it was speculated that the rhombohedral and tetragonal phases may coexist in the composition range  $x=0.10$ – $0.20$  and consequently, the existence of a temperature-independent morphotropic phase boundary (MPB). However, more structural data as a function of the temperature is still needed to assure this point and disregard other options [19]. An example of such an option is the existence of a polymorphic phase transition (PPT) and a tricritical point close to room temperature, i.e., a temperature for the coexistence of the cubic paraelectric high-temperature phase and two ferroelectric phases.

In our samples, for both  $\text{Ca}^{2+}$  compositions, the possible coexisting rhombohedral and tetragonal phases would have feeble distortions from the cubic prototype perovskite structure, featured with the merging of the (111)/(-111) and (002)/(200) doublets into single peaks around  $38^\circ$  and  $45^\circ$  in  $2\theta$ , respectively (figure 2). Such a pseudo-cubic structure is not compatible with the ferro-piezoelectric performance. This global structure observed in the ceramics here studied could, on the other hand, average a range of local polar configurations [20]. Both global and short-range structures are needed to understand the structural response. However, this is beyond the scope of this work.

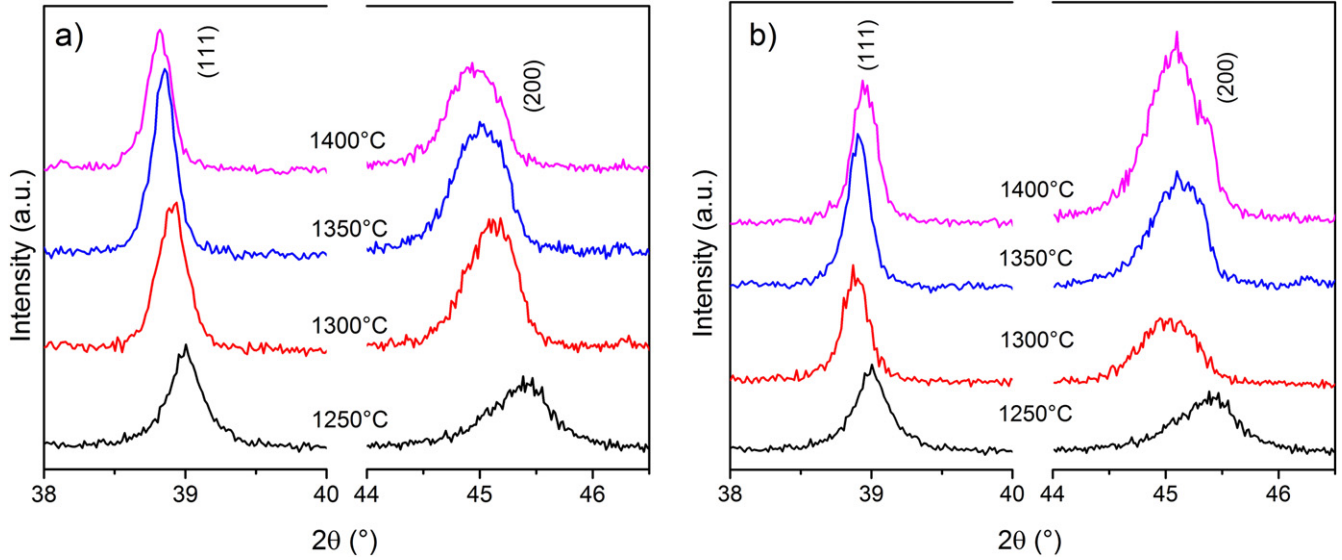
Figure 3 shows the SEM micrographs of  $x=0.10$  and  $x=0.15$  BCTZ ceramics sintered at different temperatures. As

shown, the grain size increases with the sintering temperature [10, 21] due to a higher mass transport, resulting in high densification (>90% calculated from the XRD data). Moreover, it is reported that the  $\text{Zr}^{4+}$  content helps to improve the grain growth of the BCTZ ceramics [16]. The increase in crystal size makes it easier to switch the domains, as will be discussed later. In this work the average grain size is around 0.5, 3.2, 11.1  $\mu\text{m}$  for  $x=0.10$  and 0.5, 2.1, 8.7  $\mu\text{m}$  for  $x=0.15$  at 1300, 1350 and 1400 °C, respectively.

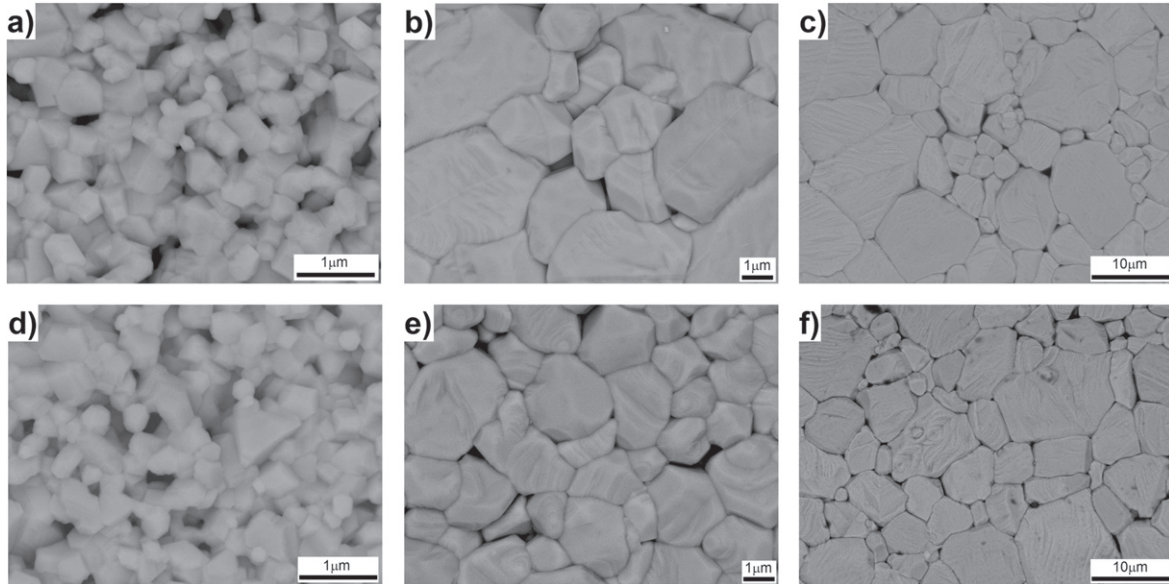
Figure 4 shows the dependence of dielectric permittivity and losses as a function of temperature for BCTZ ceramics sintered at 1400 °C. Wide permittivity peaks indicate that a diffuse phase transition is presented. In addition to the permittivity maxima related with the transition to the paraelectric phase, there is another wide transition at lower temperatures, around 40 °C. This indicates that the transition between rhombohedral and orthorhombic and between orthorhombic to tetragonal ferroelectric phases merge into a unique one [17], suggesting a coexistence of symmetries near this temperature. This effect leads to a higher number of grains being oriented favorable towards the electric field, a well-known phenomenon at the MPB.

Figure 5 displays the dielectric permittivity and losses between the two compositions under investigation. The effect that  $\text{Ca}^{2+}$  doping has over the phase transition and the Curie temperature of the ceramics is compared. In the present study, the results for dielectric permittivity correspond to previous reports [10, 22] for approximately the same grain size.

A diffuse phase transition is generally characterized by (a) broadening in the dielectric permittivity versus temperature curve; (b) a relative large separation (in temperature) between the maximum of the real dielectric constant and imaginary dielectric loss parts; (c) a deviation from the Curie–Weiss law near the maximum dielectric permittivity ( $T_m$ ); and (d) frequency dispersion of both  $\epsilon$  and  $\tan \delta$  in the transition region [23].



**Figure 2.** Magnification of the (111) and (200) peaks at the XRD patterns of BCTZ ceramics synthesized at 1250 °C and sintered at 1300, 1350, 1400 °C: (a)  $x=0.10$  and (b)  $x=0.15$ .



**Figure 3.** SEM micrographs of BCTZ sintered ceramics:  $x=0.10$  (a) 1300 °C (b) 1350 °C (c) 1400 °C;  $x=0.15$  (d) 1300 °C (e) 1350 °C (f) 1400 °C.

The dielectric permittivity of a normal ferroelectric above Curie temperature follows the Curie–Weiss law described by:

$$\frac{1}{\epsilon} = \frac{T - T_0}{C} \quad (T > T_c) \quad (1)$$

where  $T_0$  is the Curie–Weiss temperature and  $C$  is the Curie–Weiss constant. Figure 6 shows the plots of inverse dielectric permittivity of BCTZ ceramics (1400 °C) at 1 kHz versus temperature. Their respective insets correspond to the curves of  $\log(1/\epsilon - 1/\epsilon_m)$  against  $\log(T - T_m)$  at 1 kHz.

Usually, for a well-known ‘normal’ ferroelectric, the temperature of the maximum  $T_m$  corresponds to the ferroelectric–paraelectric phase transition temperature. Furthermore, the parameter  $\Delta T_m$  describes the degree of the deviation

from the Curie–Weiss law:

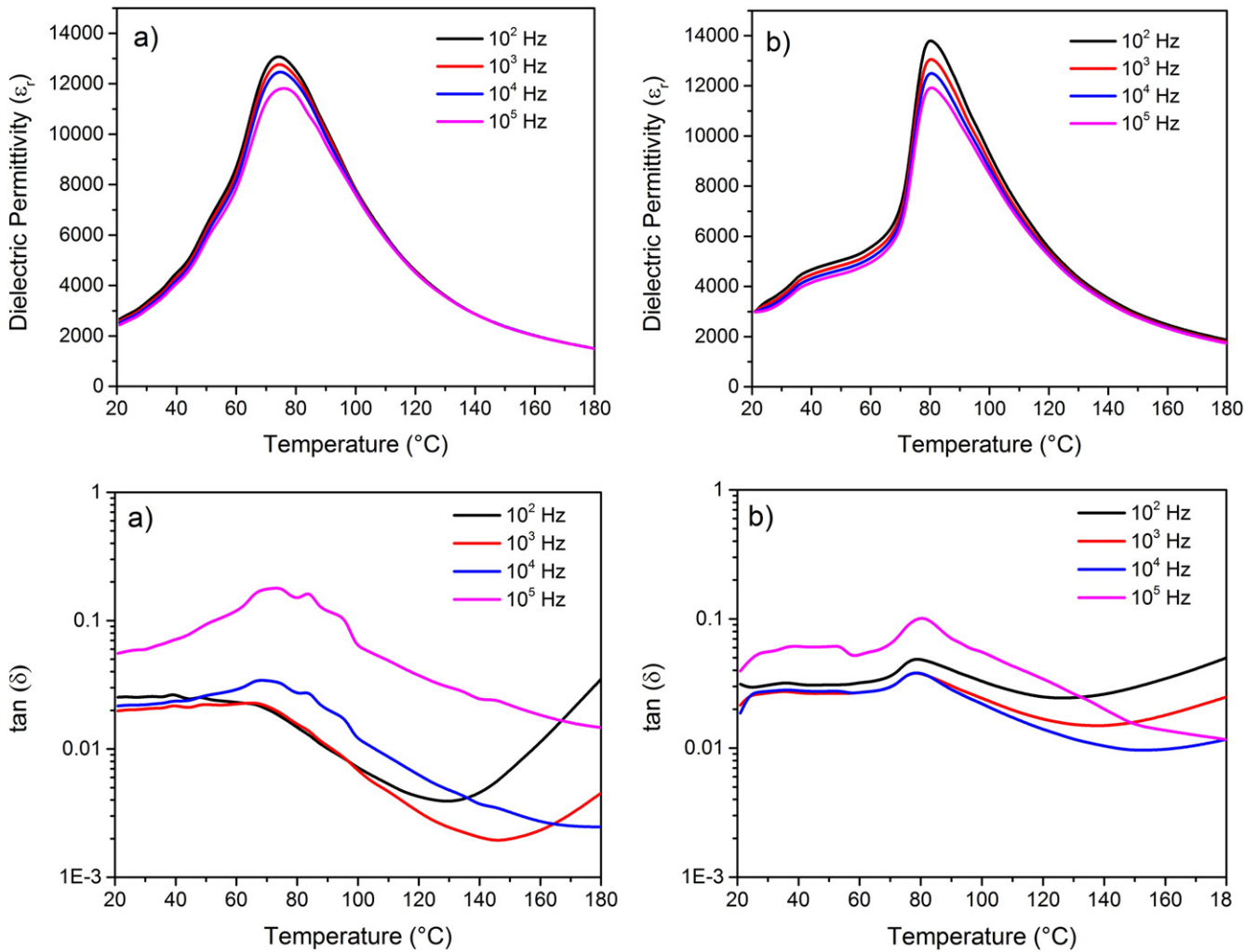
$$\Delta T_m = T_{cw} - T_m \quad (2)$$

where  $T_{cw}$  denotes the temperature from which the dielectric permittivity starts to deviate from the Curie–Weiss law and  $T_m$  represents the temperature of the dielectric maximum.

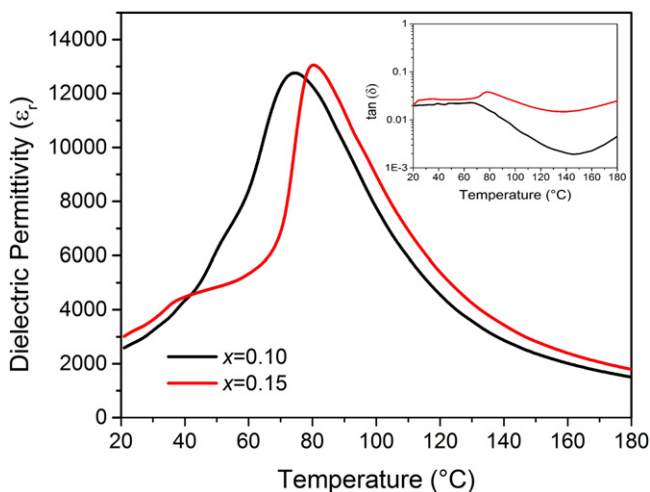
A modified Curie–Weiss expression [24] is proposed to describe the diffuseness of a phase transition:

$$\frac{1}{\epsilon} - \frac{1}{\epsilon_m} = \frac{(T - T_m)^\gamma}{C} \quad (3)$$

where  $\gamma$  (diffusion coefficient ranging) and  $C$  are assumed to be constant. The parameter  $\gamma$  gives information on the character of the phase transition:  $\gamma=1$  a normal Curie–Weiss



**Figure 4.** Dielectric permittivity and losses for BCTZ sintered ceramics (1400 °C): (a)  $x=0.10$ ; (b)  $x=0.15$ .



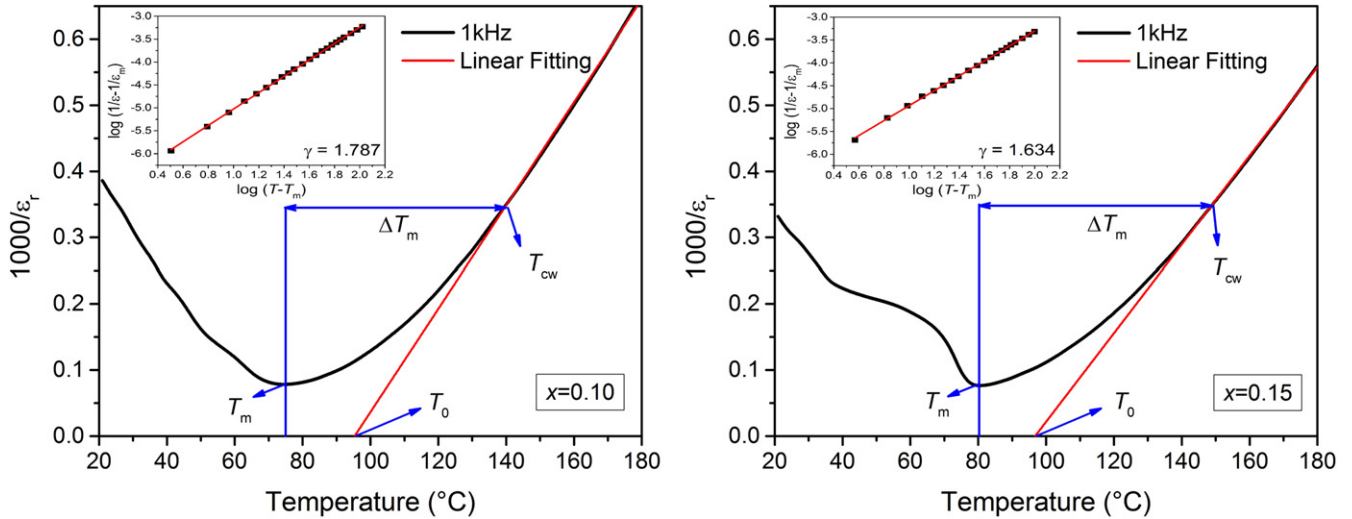
**Figure 5.** Comparative plots for dielectric permittivity and losses (inset) in BCTZ ceramics (1400 °C) at 1 kHz.

law (normal ferroelectric);  $\gamma=2$  describes a complete diffuse phase transition (ideal relaxor ferroelectric). A linear relationship for both BCTZ compositions was obtained (figure 6). The  $\gamma$  values indicate that transitions are a diffuse-like type.

Table 1 shows the  $C$  and  $\gamma$  values calculated for BCTZ ceramics in addition with a summary of the fitting parameters obtained from equations (1)–(3).

Values for  $T_0$ ,  $T_{cw}$ , and  $T_m$ , for these BCTZ ceramics are in the same trend. Moreover, the  $C$  constant follows the same tendency for various grain sizes reported with BCTZ [10]. However, the  $\gamma$  coefficient for these BCTZ studies indicates that ceramics with  $x=0.10$  compositions show a higher relaxor character.

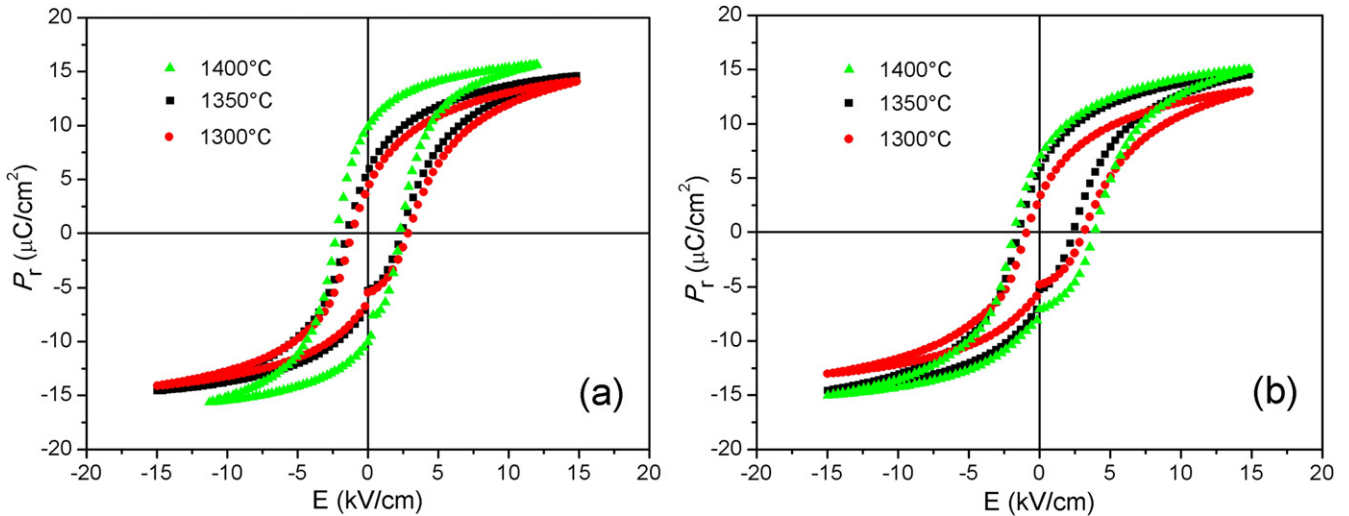
Hysteresis loops for BCTZ ceramics are shown in figure 7. Remnant polarization ( $P_r$ ) of BCTZ ceramics at  $x=0.10$  is slightly higher than for  $x=0.15$ . Furthermore, it increases with the sintering temperature. This phenomenon can be related to the easy switching of the ferroelectric domains by increasing the grain size. Table 2 shows the remnant polarization and coercive field for BCTZ ceramics sintered at different temperatures. The coercive field indicates that the ceramic is ‘soft’ in relation to the electric field because the free energy profile of the polarization rotation is anisotropically flattened at this phase boundary [8]. Furthermore, the existence of  $90^\circ$  domain walls diminishes the strain in sintered ceramics [10]. This effect also modifies the piezoelectric properties, increasing with a bigger domain size.



**Figure 6.** Inverse dielectric permittivity ( $1000/\epsilon_r$ ) as a function of temperature at 1 kHz, for BCTZ ceramics sintered at 1400 °C. Insets show the plot of  $\log(1/\epsilon - 1/\epsilon_m)$  versus  $\log(T - T_m)$  (symbols: experimental data, solid line: fitting to equation (3)).

**Table 1.** Summary of fitting parameters for dielectric behavior and anomalies in obtained BCTZ ceramics.

BCTZ	$T_0$ (°C)	$T_{cw}$ (°C)	$T_m$ (°C)	$\Delta T_m$ (°C)	$C \times 10^5$ (°C)	$\gamma$
$x=0.10$	95.3	140.2	74.8	65.4	1.539	1.787
$x=0.15$	97.2	148.9	80.3	68.6	1.799	1.634



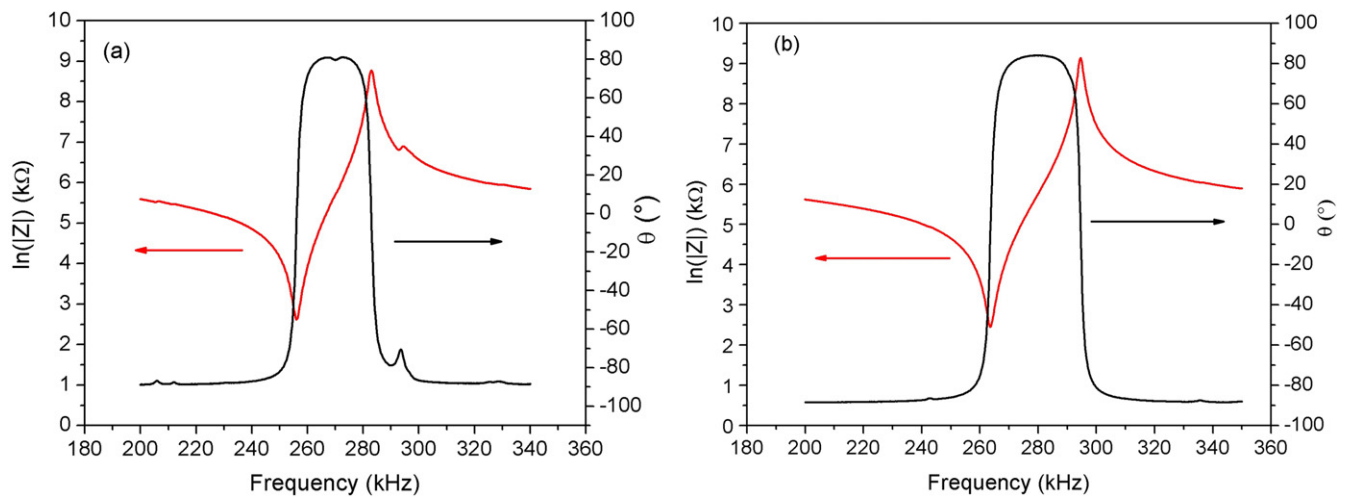
**Figure 7.** Ferroelectric loops of BCTZ ceramics at different sintering temperatures: (a)  $x=0.10$  and (b)  $x=0.15$ .

**Table 2.** Pr and EC values from obtained BCTZ ceramics.

Sample ( $x$ )	Sintering temperature °C	$P_r$ $\mu\text{C cm}^{-2}$	$E_c$ $\text{kV cm}^{-1}$
0.10	1300	4.22	2.33
	1350	5.72	2.39
	1400	9.86	2.84
0.15	1300	3.39	3.96
	1350	5.44	2.40
	1400	7.05	3.22

Table 3 summarizes the piezoelectric properties obtained for BCTZ ceramics. As seen, there exist some differences with those reported in the literature due to diverse factors [10]: intrinsic microstructural properties such as grain size, density and chemical homogeneities; measurement parameters such as frequency, maximum electric field, and temperature; and the measurement method.

Figure 8 shows the radial resonance spectra for both compositions sintered at 1400 °C for 2 h. The phase angle ( $\theta$ ) value is correlated to the  $d_{33}$ . In this work a more square-like resonance is demonstrated for the  $x=0.15$  ceramic; then the higher the  $\theta$  value, the larger the  $d_{33}$  activity.



**Figure 8.** Radial resonance spectra for BCTZ ceramics sintered at 1400 °C: (a)  $x=0.10$  and (b)  $x=0.15$ .

**Table 3.** Summary of piezoelectric properties from obtained BCTZ ceramics.

Sample	Sintering temperature	$\rho$	$d_{33}$	$-d_{31}$	$k_p$	$s_{11}^E$	$-s_{12}^E$	$g_{33}$	$-g_{31}$
(x)	°C	%	$10^{-12} \text{ C N}^{-1}$	$10^{-12} \text{ C N}^{-1}$	%	$10^{-12} \text{ m}^2 \text{ N}^{-1}$	$10^{-12} \text{ m}^2 \text{ N}^{-1}$	$10^{-3} \text{ Vm N}^{-1}$	$10^{-3} \text{ Vm N}^{-1}$
0.10	1300	94	175	69	22	11.8	4.3	6.6	2.6
	1350	95	285	96	31	11.3	3.4	11.7	4.0
	1400	97	300	150	48	11.0	3.3	11.9	5.3
0.15	1300	96	210	78	27	14.1	7.4	8.4	3.1
	1350	97	330	132	44	11.7	4.3	13.5	5.4
	1400	98	410	154	50	10.8	3.1	16.3	6.1

Although there are small differences in density among the samples sintered at different temperatures for both compositions, the piezoelectric performance drastically changes. Making a comparison between samples sintered at 1400 °C, the  $d_{33}$  value for  $x=0.15$  is higher than the value obtained for  $x=0.10$ , despite both showing a square-like shape of radial resonance. Moreover, as temperature increases for  $x=0.15$  the  $d_{33}$  value is almost doubled. The same trend is observed for  $d_{31}$  and  $k_p$ . As mentioned earlier, it seems that the grain size has more influence over the piezoelectric properties than the density, since all the obtained ceramics are dense enough (theoretical density for  $x=0.10$ :  $5.83 \text{ g cm}^{-3}$  and  $x=0.15$ :  $5.75 \text{ g cm}^{-3}$ ). This effect could be assigned to the contribution of  $90^\circ$  ferroelectric domains in ceramics with large crystals [10], as it enables the switching/rotation of the polarization vector.

#### 4. Conclusions

BCTZ ceramics were obtained by solid-state synthesis using a reduced thermal treatment at 1250 °C (2 h) and different sintering temperatures (1300, 1350, 1400 °C for 2 h). Two compositions were studied ( $x=0.10$ ,  $x=0.15$ ). Both systems show a single perovskite phase type structure observed by x-ray diffraction analysis. It was found that in both compositions, ceramics sintered at 1400 °C exhibit excellent piezoelectric and ferroelectric properties due to an average increase

in grain size, which promotes the ceramic polarization. From our results, it is also worth noting the production of high-sensitivity ( $d_{33} > 250 \text{ pC/N}$ ) piezoelectric ceramic bodies of the two compositions using a conventional solid-state synthesis method and temperature times as low as 1350 °C for 2 h with dense microstructure and average grain size in the  $1 \mu\text{m}$  range. These BCTZ ceramics are promising candidates for many electronic micro-devices.

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