



A stable multiply twinned decahedral gold nanoparticle with a barrel-like shape



Ulises Santiago ^a, J. Jesus Velázquez-Salazar ^a, John Eder Sanchez ^a, Francisco Ruiz-Zepeda ^{a,1}, José Eduardo Ortega ^a, José Reyes-Gasga ^b, Lourdes Bazán-Díaz ^{a,b}, Israel Betancourt ^{a,2}, Edgar F. Rauch ^{c,d}, Muriel Veron ^c, Arturo Ponce ^a, Miguel José-Yacamán ^{a,*}

^a Department of Physics and Astronomy, University of Texas at San Antonio, One UTSA Circle, San Antonio, TX 78249, USA

^b Instituto de Física, UNAM, Circuito de la Investigación s/n, Col. Ciudad Universitaria, Coyoacán, México D. F. 04510, Mexico

^c Univ. Grenoble Alpes, SIMAP, F-38000 Grenoble, France

^d CNRS, SIMAP, F-38000 Grenoble, France

ARTICLE INFO

Article history:

Received 9 July 2015

Accepted 14 September 2015

Available online 25 September 2015

Keywords:

Decahedral barrel-like nanoparticle

Coherent electron diffraction

Crystal orientation maps

Precession electron diffraction

ABSTRACT

In the present work we report a modified configuration of a multiply twinned decahedral barrel-like nanoparticle. The nanoparticle is stabilized by a selective chemical etching, which yielded a truncated shape with multiple high-index planes on the surface. The surface planes and the shape of the nanoparticle have been characterized using electron microscopy techniques, including scanning and transmission electron microscopies and electron diffraction. Coherent electron diffraction revealed streaked Bragg reflections which are in agreement with the high-index facets measured from electron microscope projected images. Crystal orientation mapping assisted by precession electron diffraction was performed under nanobeam conditions giving information about the structure from different orientations. The combination of these techniques allowed us to resolve the 3D structure of the modified nanoparticle and dismiss any possible ambiguities from the interpretation of the projected images in direct space.

© 2015 Elsevier B.V. All rights reserved.

1. Introduction

In the present state of the art in nanotechnology the full control of shape, size and crystal structure of nanoparticles is a fundamental aspect to develop practical applications. In order to achieve the desired structure with also high yield and reproducibility is necessary to fully understand the factors that define the growth of nanoparticles. The growth kinetics results in different shapes, in this way there is an interplay between particle total energy and the kinetics of the growth. Their stability depends mainly upon the size range desired. Different shapes can be produced during growth by balancing the energies binding within the nanoparticle and the ones at the surface planes, where the free energy plays a crucial role in the stabilization of the nanoparticles, and hence giving them their final shape [1,2]. In the smallest size range, nanoparticles and clusters are stabilized by the use of passivating ligands depending on the number of atoms. The diameter of these nanoparticles ranges from a few atoms up to a few nanometers [3–5]. The

next size range includes tens of nanometers in diameter. In this range, plasmonic effects start to appear and the shape can be mostly controlled by kinetics. A prime example being the nanorods, nanowires and hollow particles [6,7]. The next size interval includes larger particles (>100 nm). In this case the shapes are mostly polyhedral nanoparticles based on the Wulff construction formalism [6,8], which has been adapted to multiply twinned particles and oxide nanoparticles [9,10] with recent modifications by Ringe et al. [11,12]. In addition, other reports also consider the edge energies obtained by atomistic calculations to determine and predict some shapes as a function of the number of atoms distributed along the edges [13,14]. The way particles become stable is by faceting, allowing the formation of structures such as decahedral or icosahedral structures even on this size interval [15,16]. In the present work we investigate the shape of a modified truncated decahedron which exhibits a barrel-like shape with multiple facets using single particle diffraction and precession electron diffraction. This serves as a prime example of the very complex structures that can be produced by faceting.

2. Methods

The nanoparticles have been prepared by a controlled chemical synthesis following the polyol method. Initially, an ethylene glycol (EG) solution containing polyvinylpyrrolidone (PVP) (0.209 g, $M_w = 55,000$)

* Corresponding author. Tel.: +1 210 458 6954.

E-mail address: miguel.yacamán@utsa.edu (M. José-Yacamán).

¹ Present address: National Institute of Chemistry, Hajdrihova 19, 1000 Ljubljana, Slovenia.

² On Sabbatical leave from Instituto de Investigaciones en Materiales, Universidad Nacional Autónoma de México, C.P., D.F. 04510 México.

was heated to boiling, into which an AgNO_3 (0.138 ml, 0.6 mM) solution in EG was added rapidly under magnetic stirring. After 3 min, aliquots of a $\text{HAuCl}_4 \cdot 3\text{H}_2\text{O}$ (0.162 ml, 2.5 mM) solution were added each 7 s drop by drop. The solution turned to red and finally changed to reddish brown, indicating precipitation of gold. Then, the mixture (growth solution) was stirred for 10 min more and it was quickly cooled to room temperature. The colloidal solution was washed with ethanol and acetone several times to remove excess PVP and EG: 1 ml of gold nanoparticles was added to 11 ml of acetone and the nanoparticles were collected by centrifugation at 6000 rpm. Products were dispersed in 3 ml of ethanol and 8 ml of acetone and precipitated one more time by centrifugation. This process is repeated several times. The final product of gold nanoparticles was dispersed in ethanol. Our method produces different truncated nanoparticles such as decahedra, cuboctahedra, plates and nanorods all of them produce facets with low-energy $\{111\}$ surfaces, which are produced by the etching during the chemical synthesis and are in agreement with the surface faceting reported in our previous work [14]. The truncated decahedron shape is herein analyzed. An example of the obtained nanoparticles with different shapes is shown in Figure S1 of the supplemental material.

A drop of the colloidal solution was deposited onto copper holey carbon grids for characterization. The analysis of the shape and projected planes of the nanoparticles was performed by scanning electron microscopy (SEM) and transmission electron microscopy (TEM). The SEM images were obtained using a field emission gun scanning transmission electron microscope Hitachi S5500 operated at 30 kV. The TEM imaging information was obtained using field emission gun JEOL microscopes, 2010F and 200F ARM both of them operated at 200 kV. In the present work we determine the geometry of modified gold decahedra nanoparticles by using selected area electron diffraction (SAD), precession electron diffraction-assisted automated crystal orientation mapping (PED ACOM-TEM) [17] and coherent electron diffraction (CED) [18,19]. These electron diffraction techniques provided useful information about the crystalline structure and surface plane facets of the barrel-like nanoparticle. To perform PED and CED, a small area is irradiated under nanobeam electron diffraction conditions, using the smallest condenser aperture and reducing the diameter of the electron irradiation to a few nanometers. In PED we used a probe size of around 1.5 nm in diameter, while in CED, in order to illuminate the whole particle, the probe size was slightly larger than the nanoparticle size (more than 100 nm). A CED pattern shows streak lines in the Bragg reflections; these streak lines are measured with respect to the Bragg reflection and are related to high index crystallographic planes showing directions normal to the facets of the nanoparticle. These streaks come from the intersection of the Ewald sphere with reciprocal lattice rods

of the nanoparticle facets. The intersection is almost tangential when the facet is normal to the incident electron beam. The sharper the facet the longer the length of the streaks [19]. On the other hand, PED ACOM-TEM gives the full crystallographic information of the object at a specific area of the region of interest (ROI), by providing a crystallographic map at the nanometer scale and making easier the identification of crystalline phases in semiconductor materials, metallurgy and nanostructured materials [20–22]. In the ACOM-TEM technique the electron beam is scanned across the sample and collects the electron diffraction patterns using an ultrafast charge-coupled device (CCD) camera attached to the viewing screen of the microscope. All of the electron diffraction patterns collected are then indexed using an automatic matching procedure, identified by a cross-correlation with preselected theoretical templates from a database. The ACOM-TEM technique is assisted by PED to enhance the quality of the diffraction pattern matching, since PED reduces the dynamic effects arising in electron diffraction [23]. PED experiments were conducted using Nanomegas hardware and software package [24]; further details of the experimental set-up and from the analysis process are available in the supplementary material.

3. Results and discussion

The truncated decahedral Au nanoparticles exhibit a barrel-like form where the base is pentagonal in shape (similar to a truncated decahedral particle) and the lateral view is formed by 50 planes faceted one to each other $5(\text{sides}) \times 2(\text{bi-pyramid})(2 \times \{320\} + 2 \times \{221\} + 1 \times \{111\})$. Fig. 1 shows the scanning electron microscopy images of the nanoparticle. These facets stand the body at an angle of 143° to the base; the identified planes are labeled in Fig. 1a and b. Table 1 presents the theoretical $\{hkl\}$ planes and the measured angles from the SEM images of Fig. 1. It is worth noting that in the drawing of Table 1 the particle is seen along the direction $[110]$, and that, e.g. the $\{320\}$ and $\{221\}$ facets are not parallel to $[110]$; in fact the $\{221\}$ facet is 13° of inclination with respect to $[110]$ and about the same goes with $\{320\}$. However, despite this inclination, the angle between the facets is practically the same. That is, the angles in Table 1 could vary 1 or 2° from the observed value, but they will be very close to the calculated value.

TEM imaging has been used to confirm the angles measured in SEM images. Within the nanoparticle the atomic distribution corresponds to a regular multiple twin decahedral structure. In Fig. 2, TEM images are shown in two different orientations: along the $\langle 110 \rangle$ direction and perpendicular to the $\langle 110 \rangle$ (the side view direction). The fcc unit cell of Au has been taken as the reference system. The side view in Fig. 2a shows the projected planes $\{320\}$ and $\{111\}$ as well as the truncated plane of

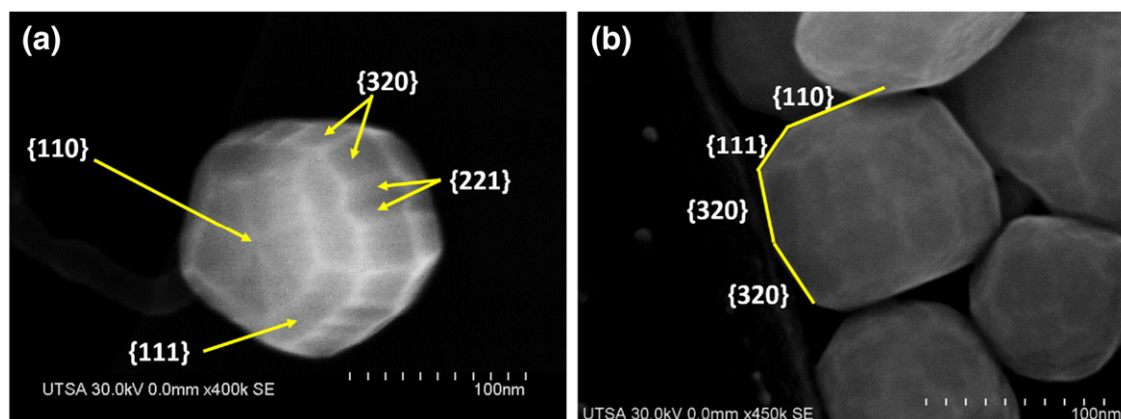
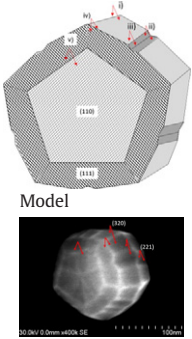


Fig. 1. SEM images of the nanoparticle. a) Front view; b) side view. The arrows in (a) indicate the surface planes of the surface nanoparticle. The yellow lines in (b) indicate the edge surface.

Table 1
Interplanar angles measured for the barrel-like nanoparticle and compared with their theoretical values.

Projection	Interplanar angles	Planes $\{h_1k_1l_1\}/\{h_2k_2l_2\}$	Angles (in degrees)	
			Theoretical	Experimental
 <p>Model SEM micrograph from Fig. 1a</p>	i)	$\{3\ 2\ 0\}/\{3\ 2\ 0\}$	133.81	135 ± 3
	ii)	$\{2\ 2\ 1\}/\{2\ 2\ 1\}$	152.73	155 ± 3
	iii)	$\{3\ 2\ 0\}/\{2\ 2\ 1\}$	157.59	156 ± 2
	iv)	$\{3\ 2\ 0\}/\{1\ 1\ 1\}$	143.19	140 ± 3
	v)	$\{1\ 1\ 0\}/\{1\ 1\ 1\}$	144.74	144 ± 2

the decahedron $\{110\}$. Angles are consistent with those measured in Fig. 1; in Fig. 2d and e we have measured the angles on the particle oriented close to the $\langle 110 \rangle$ zone axis. Especially, Fig. 2e shows the contours of the nanoparticle pictured in Fig. 2d in order to measure more precisely the angles of the facets in the nanoparticle. In Fig. 2b, the corresponding SAD pattern from the side view of the particle (Fig. 2a) is shown.

The spots in the SAD pattern have been identified as a combination of two overlapped zone axes $[001] + [112]$. The overlapping of these two zone axes is simulated in Fig. 2c and correlates well with the reflections observed in the experimental SAD pattern of Fig. 2b. The extra spots that appear in Fig. 2b arise from the combination of double

diffraction due to the multiple twinned structure and the nano size of the particle [7,25]. It is well known that the pattern in Fig. 2f is the result of the convolution of five $[110]$ patterns each one rotated by 72° . The $\{221\}$ planes projected from the side view are not totally symmetric in the five sides of the decahedron due to a slight tilting of the particle out of the $\langle 110 \rangle$ twinning axis.

Once the relationship between the axis $[110]$ and the angle is known in the truncated decahedral nanoparticle, it is possible to know which planes are observed when the particle is seen side view. In the cubic system, the direction $[hkl]$ is perpendicular to the plane (hkl) and given that the observation axis is along one of the $\{110\}$ axes family, the diffraction pattern along the $[110]$ zone axis is used to calculate these planes. Using

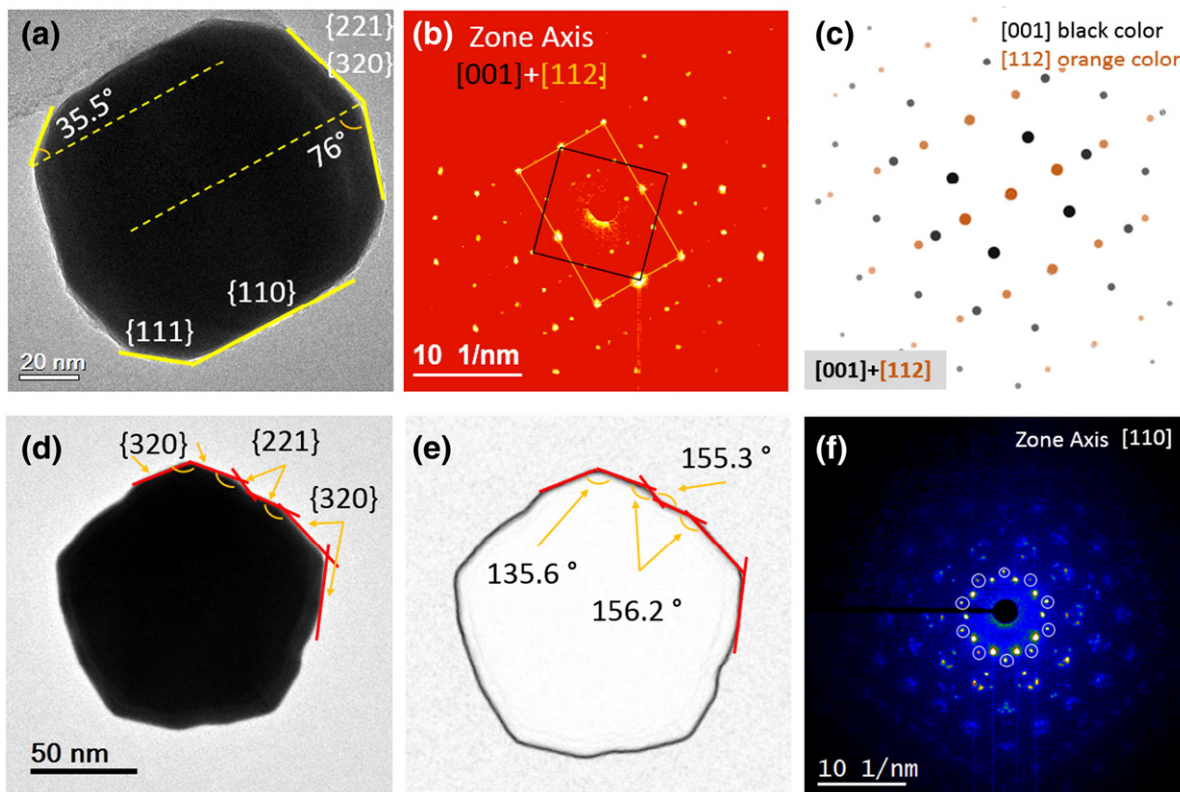


Fig. 2. TEM images from (a) side view, (d) close to the five-fold (along $\langle 110 \rangle$ direction) and (e) contours filtered from (d). Projected planes and angles are measured from both orientations. SAD patterns in (b) and (f) are recorded along the side view and close to the five-fold, respectively. The circled diffraction spots in (f) correspond to the $\{002\}$ planes. (c) Simulation of two overlapped electron diffraction patterns $[001] + [112]$.

solely SAD, the high index faceting cannot be identified with ease; in such cases, CED has demonstrated to be a powerful tool for the 3D interpenetration of high index facets in polyhedral nanoparticles and other nanostructures with atomic steps on their surfaces [18,19]. Following this line of thought we have performed CED in two directions as shown in Fig. 3. The CED pattern in Fig. 3a shows an inset of the illumination conditions for the CED method along the five-fold direction. A similar image is shown in Fig. 3b for the nanoparticle oriented slightly out of the five-fold direction. In both cases we have observed streak lines which the angles measured have revealed facets also identified in images from Figs. 1 and 2. For instance, the CED pattern in Fig. 3a shows streak lines oriented at 54° with respect to the $\langle 111 \rangle$ Bragg reflection. If we consider the particle oriented along $\langle 110 \rangle$, and taking the cubic system as the reference, the plane $\{111\}$ is inclined more than the planes $\{320\}$ and $\{221\}$ from the side part of the nanoparticle. Directions $[221]$ and $[320]$, which are normal to the planes $\{221\}$ and $\{320\}$ respectively, made an angle of about 13° with the $[110]$ axis, approximately, while the angle between planes $\{111\}$ and $\{110\}$, or directions $[111]$ and $[110]$, is 32° , approximately. This is what is observed in the produced streak lines. This 54° multiplied 10 times gives the sum of the inner angles in a pentagon. The pattern in Fig. 3b corresponds mainly to the $[111]$ zone axis with some contributions of the $[112]$ zone axis

which produces the streak lines for the faces $\{320\}$ and $\{221\}$. This was also confirmed by electron tomography from which we performed 3D reconstruction and visualization using the algebraic reconstruction technique (ART) [26]. The details about the electron tomography and 3D reconstruction are available in the supplementary material as well as the video of the 3D reconstruction.

The analysis of the geometry of these nanoparticles has been performed using the PED ACOM-TEM technique. A full set of diffraction patterns was collected while the precessing electron beam was scanning the ROI. The step size used in the scanning was 1.5 nm, about the same as the probe size. The results of the indexation are later displayed as a crystal orientation map of the ROI. In Fig. 4, the barrel-like nanoparticle has been oriented along the five-fold symmetry. Each orientation (colored) map, refers to a particular viewing direction, i.e. the color is assigned based on the crystal orientation along this particular direction. Experimental details about the PED ACOM-TEM measurements are available in the supplementary material. Fig. 4 shows also the color code map for the crystal directions (Fig. 4b) and a model of the five-fold nanoparticle orientation, the axes and e-beam directions (Fig. 4a). In Figure S3 of the supplementary material an example of the matching between a particular pattern and the relevant template is shown. Fig. 4c, d and e shows the orientation of each part of the decahedron observed from the three main reference directions x, y, and z, respectively. Fig. 4c shows the planes $\{111\}$ (blue color) and $\{110\}$ (green color) which are consistent with the view using as a reference the “x” axis. Note that in Fig. 4d all crystals show the same color, this is clearly expected because the orientation of the nanoparticle is such that the viewing direction (e-beam direction) is parallel to the common $[110]$ direction. The colors displayed in the orientation maps of Fig. 4, correspond to the $[111]$, $[001]$ and $[101]$ reference directions and to the $[112]$, $[113]$ and $[212]$ directions, having all crystals an orientation of 72° between them.

In addition, we have recorded crystal orientation maps on two different tilt directions, 50° away from each other. First, the particle analyzed was tilted close to a side view at $+25^\circ$ from a starting point. In this orientation (Fig. 5 – top set) the five crystals are identified and indicated by different colors in at least two of the orientation maps, but the interfaces are very poorly defined. This is because the common $[110]$ axis is inclined 60° with respect to the beam direction (z-direction). Consequently, crystals are massively overlapping in the thickest part of the nanoparticle and the selection of the predominant diffraction pattern is confused by the strong signal related to the additional patterns. However, the twin relationships have been confirmed in all cases by computing the disorientation between two adjacent crystals (60°) and the associated rotation axis ($[111]$). A second scan was performed after tilting the sample 50° around the sample holder axis (Fig. 5 – bottom set). In this new orientation, the common $[110]$ axis is far closer from the incident beam direction (20°) and overlapping artifacts are reduced. In particular, there is a substantial improvement in crystal interface localization. Again, the five crystals are recognized and the twin characters confirmed. Note that for Fig. 5 the reference axes have been tilted 44° clockwise. This is in order to represent the orientation in the physical reference frame related to the sample holder. Obviously, the crystallographic orientations along x, that is, along the sample holder main axis, should not be altered with the tilt angle. Indeed, the colors in the orientation map labeled ‘x’ remain unchanged for every crystal before and after tilting (compare Fig. 5c and c’) leading to an easier interpretation of the results.

To gain a better view of the 3D configuration of the nanocrystal, correlation contours are added in Fig. 5f and f’. These contours correspond to the average value of the cross-correlation indexes between a given diffraction pattern and the surrounding ones. A large value (white) is obtained when the diffraction patterns are similar, i.e. within a crystal. The value decreases if a change in the diffracting signal is noticed. Dark lines are expected to highlight grain or phase boundaries. The contrast is enhanced when the boundaries are aligned with the beam direction (abrupt change of diffraction patterns). In the present case, the

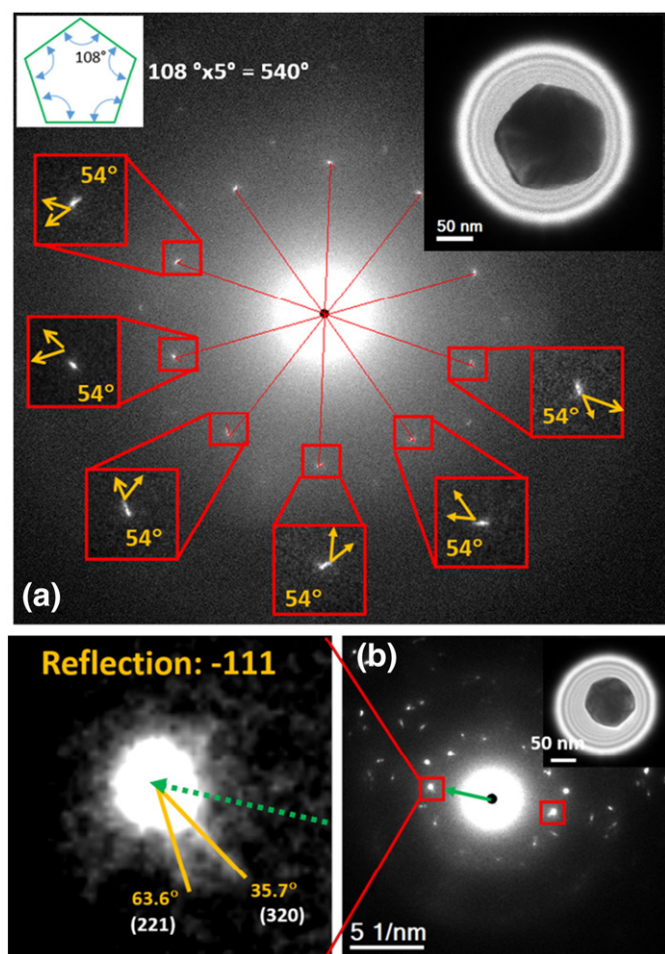


Fig. 3. Coherent electron diffraction patterns from the truncated decahedral particle. (a) Along plan-view; (b) along side-view. The pattern in (a) is the result of the convolution of five $[110]$ patterns rotated by 72° each grain in the nanoparticle. The yellow arrows correspond to the direction of reference, taken from the transmitted beam, and to the direction of the streak, yielding 54° between them in all cases. The pattern in (b) corresponds to the $[111]$ zone axis together with some spots of the $[112]$ zone axis electron diffraction pattern.

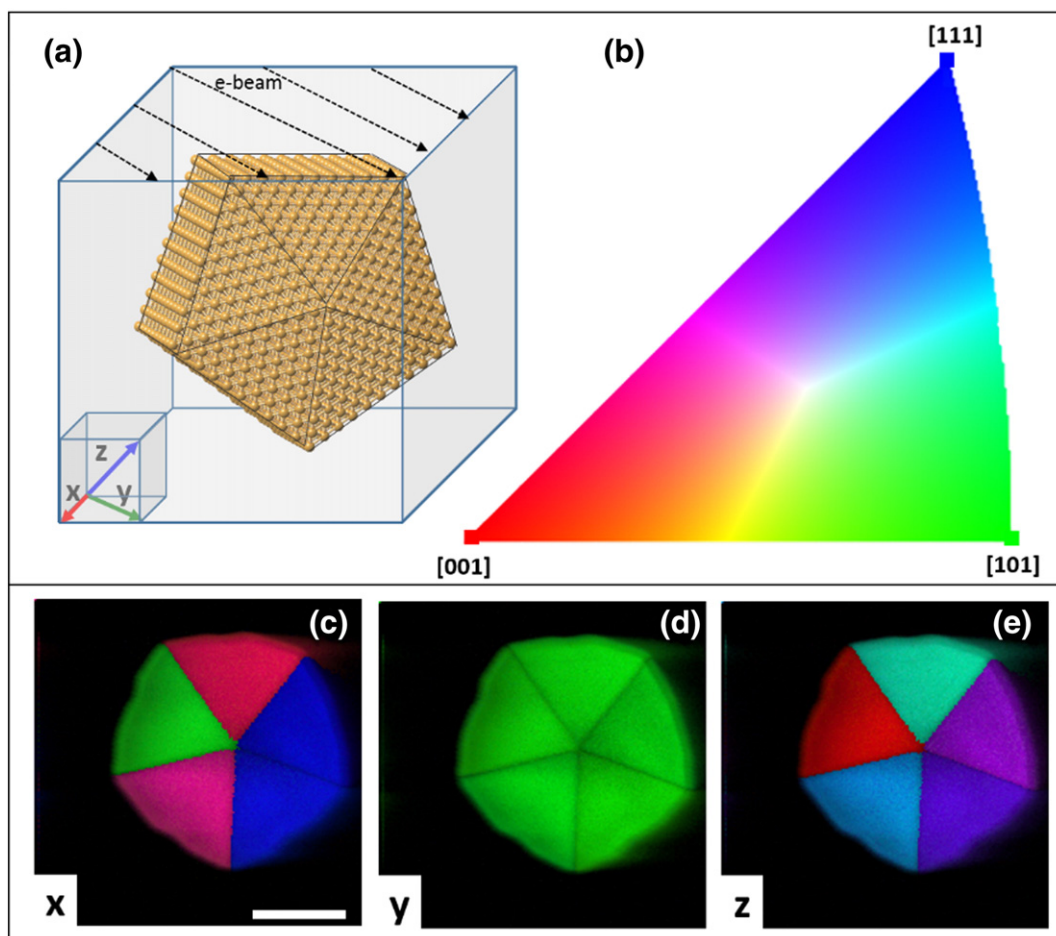


Fig. 4. (a) Decahedron model and the frame of reference for the crystal orientation maps. (b) Color code inverse pole figure taken from the inverse pole diagram obtained along the [001] direction (included in Figure S3). Crystal orientation maps corresponding to x, y and z directions, (c), (d) and (e) respectively (scale bar is 50 nm).

inner twin boundaries are detected, although not perfectly but sufficiently to guess the orientations of the crystals before and after the tilt. From these views there is clearly no twin cutting the nanoparticle in cross section along the direction [110]. Additional particles with and without multiple twins were tested with the same method and for the case of twinned particles the contours are clearly visible. These images are available in the supplementary information Figure S5.

4. Conclusions

We have here reported a modified stable decahedral truncated nanoparticle with a “barrel-like” shape. Electron diffraction patterns revealed the facets on its surface and orientation phase maps were used to distinguish individual units within the nanoparticle. The particle has an equivalent pentagonal symmetry with respect to known nanoparticles.

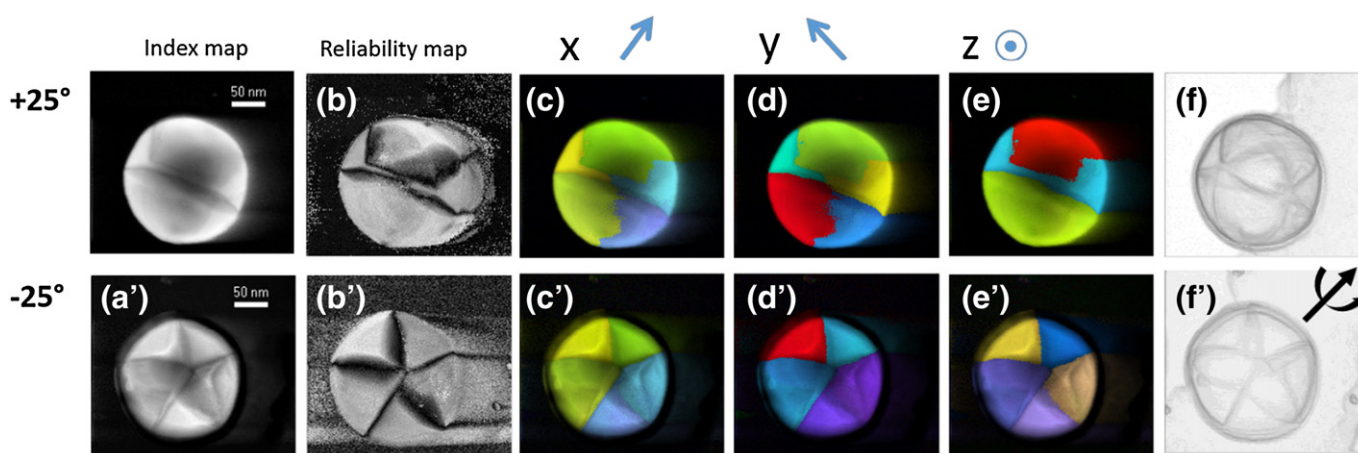


Fig. 5. Visualization maps for index, reliability, crystal orientation with reference directions x, y and z, and contour representation, obtained for a decahedral barrel-like nanoparticle at two different tilts $+25^\circ$ and -25° with respect to a starting point.

The planes of the nanoparticle have been identified by crystal orientation maps using ultra-fast scanning electron diffraction under precession mode and coherent electron diffraction, both diffraction modes not only complement each other, but validate and give consistency to the observed geometry.

Acknowledgments

The authors would like to acknowledge the NSF PREM DMR #0934218, the Department of Defense #64756-RT-REP and CONACYT #250836. I. Betancourt acknowledges the financial support from PASPA-DGAPA, UNAM and CONACYT, Mexico, during his sabbatical leave. E.F. Rauch acknowledges Ákos K. Kiss from the University of Pannonia – Hungary for suggesting the correlation contour principle prior to the developments shown in the paper. The microscopy work was supported by the National Institute on Minority Health and Health Disparities (NIMHD) in the program Research Centers in Minority Institutions Program (RCMI) Nanotechnology and Human Health Core (G12MD007591).

Appendix A. Supplementary data

Supplementary data to this article can be found online at <http://dx.doi.org/10.1016/j.susc.2015.09.015>.

References

- [1] A.S. Barnard, *Nanoscale* 6 (2014) 9983.
- [2] F. Baletto, R. Ferrando, *Rev. Mod. Phys.* 77 (2005) 371.
- [3] M. José Yacamán, J.A. Ascencio, H.B. Liu, J. Gardea-Torresdey, *J. Vac. Sci. Technol. B* 19 (2001) 1091.
- [4] A. Dass, *J. Am. Chem. Soc.* 133 (2011) 19259.
- [5] D. Bahena, N.U. Santiago, A. Tlahuice, A. Ponce, S.B.H. Bach, B. Yoon, R.L. Whetten, U. Landman, M. Jose-Yacamán, *J. Phys. Chem. Lett.* 4 (2013) 975.
- [6] S. Ino, *J. Phys. Soc. Jpn.* 27 (1969) 941.
- [7] J. Reyes-Gasga, J.L. Elichiguerra, C. Liu, A. Camacho-Bragado, J.M. Montejano-Carrizales, M. Jose-Yacamán, *J. Cryst. Growth* 286 (2006) 162.
- [8] G. Wulff, *Kristallogr. Miner.* 34 (1901) 449.
- [9] L.D. Marks, *Philos. Mag.* A 49 (1984) 81.
- [10] J.A. Enterkin, K.R. Poeppelmeier, L.D. Marks, *Nano Lett.* 11 (2011) 993.
- [11] E. Ringe, R.P. Van Duyne, L.D. Marks, *J. Phys. Chem. C* 117 (2013) 15859.
- [12] E. Ringe, R.P. Van Duyne, L.D. Marks, *Nano Lett.* 11 (2011) 3399.
- [13] J.C. Hamilton, *Phys. Rev. B* 73 (2006) 125447.
- [14] G. Casillas, J.J. Velazquez-Salazar, M. Jose-Yacamán, *J. Phys. Chem. C* 116 (2012) 8844.
- [15] E. Carbo-Argibay, B. Rodriguez-Gonzalez, I. Pastoriza-Santos, J. Perez-Juste, L.M. Liz-Marzan, *Nanoscale* 2 (2010) 2377.
- [16] Z. Zhou, N. Tian, Z. Huang, D. Chen, S. Sun, *Faraday Discuss.* 140 (2008) 81.
- [17] E.F. Rauch, M. Véron, *Mater. Charact.* 98 (2014) 1.
- [18] J.M. Zuo, M. Gao, J. Tao, B. Li, R. Twisten, I. Petrov, *Microsc. Res. Tech.* 64 (2004) 347.
- [19] A.B. Shah, S.T. Sivapalan, B.M. DeVetter, T.K. Yang, J. Wen, R. Bhargava, C.J. Murphy, J.M. Zuo, *Nano Lett.* 13 (2013) 1840.
- [20] I. Ghamarian, Y. Liu, P. Samimi, P.C. Collins, *Acta Mater.* 79 (2014) 203.
- [21] A. Avilov, K. Kuligin, S. Nicolopoulos, M. Nickolskiy, K. Boulahya, J. Portillo, G. Lepeshov, B. Sobolev, J.P. Collette, N. Martin, A.C. Robins, P. Fischione, *Ultramicroscopy* 107 (2007) 431.
- [22] F. Ruiz-Zepeda, Y.L. Casallas-Moreno, J. Cantu-Valle, D. Alducin, U. Santiago, M. Jose-Yacamán, M. Lopez-Lopez, A. Ponce, *Microsc. Res. Tech.* 77 (2014) 980.
- [23] E.F. Rauch, M. Véron, J. Portillo, D. Bultreys, Y. Maniette, S. Nicolopoulos, *Microsc. Anal.* 22 (2009) S5.
- [24] <http://www.nanomegas.com/> (accessed: 7, 2015).
- [25] J.L. Elichiguerra, J. Reyes-Gasga, M. Jose Yacamán, *J. Mater. Chem.* 16 (2006) 3906.
- [26] C. Messaoudi, T. Boudier, C.O. Sanchez-Sorzano, S. Marco, *BMC Bioinf.* 8 (2007) 288.