



Optical properties of nanocrystalline La₂O₃ dielectric films deposited by radio frequency magnetron sputtering



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ABSTRACT

La₂O₃ thin films were successfully synthesized by r. f. magnetron sputtering technique. The effect of power, deposition time and substrate temperature on the formation and optical properties of the films was investigated. X-ray diffraction (XRD) studies revealed the formation of hexagonal phased La₂O₃ thin films. The influence of sputtering parameters on chemical composition and surface species was studied by X-ray photoelectron spectroscopy (XPS). The optical properties were investigated in the wavelength range of 200–1100 nm. The samples were modelled as a three-phase optical model. Optical constants were calculated at 2 eV from classical dispersion model based on the single Lorentz for dielectric materials.

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1. Introduction

Actually, a lot of research efforts is devoted to search for new alternative dielectric materials for complementary metal-oxide-semiconductor (CMOS) devices and materials that increase the optical properties [1–10]. CMOS technology is the basic semiconductor technology for microprocessors, memories and application specific integrated circuits, whereas optical structures are commonly used as non-volatile memory pyro-electric detectors and microwave devices.

Particularly, transparent thin films with good optical properties are demanded in integrated optics. In this relation, lanthanum oxide possesses a lot of properties attractive for industrial and technological applications [11–14]. La₂O₃ films are conventionally prepared using various deposition techniques such as chemical vapor deposition, spray pyrolysis [15–17], electron-beam evaporation, pulsed-laser evaporation and vacuum evaporation [18–20]. Besides, the reactive dc magnetron sputtering method is very promising for the La₂O₃ film fabrication in an amorphous state [21]. A route very few explored to produce this kind of rare earth thin films is the radio frequency (r. f.) magnetron sputtering technique which is one of the most widely used in the manufacture of integrated circuits, resistant materials delaying

the corrosion process and optical devices [22–27]. However, the dependence of particle size and optical properties from the synthesis parameters to obtain nanocrystalline lanthanum oxide has not been yet reported. The present study is aimed to find the conditions for the successful deposition of nanocrystalline La₂O₃ thin films on Si (100) substrates by the r. f. magnetron sputtering technique and to assess the differences in the structure, morphology, composition and optical properties depending on deposition parameters.

2. Experimental details

Lanthanum oxide films were deposited on Si (100) substrates using r. f. magnetron sputtering apparatus from a La₂O₃ target (99.99% purity with copper bonding, Plasmaterials) in Ar atmosphere of ultra-high purity. Prior to the La₂O₃ deposition, the Si substrates were cleaned in sonication bath for 20 min with water-isopropyl alcohol-water, thereafter samples were blown with compressed air and dried at room temperature. The base pressure of the sputter chamber was about 2.66 Pa (20 mtorr) with a substrate-target distance of 0.06 m. From previous studies and experience of the team group, it was decided to adjust the r. f. sputtering power at 60 and 90 W; with an Ar flow of 30 sccm. The synthesis was realized at each sputtered power using different times (25, 40 and 60 min) and the substrate temperature adjusted previous to the deposition was 200 °C.

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2.1. Chemical, morphological and structural characterization

The chemical composition and band structure of La_2O_3 films were characterized by x-ray photoelectron spectroscopy (XPS) using a VG Microtech Multilab ESCA 2000 system with a CLAM MCD detector, Al K_{α} ($h = 1453.6$ eV) radiation, operating at 8×10^{-7} Pa, a pass energy of 20 eV and an energy resolution of 0.1 eV. To obtain the elemental chemical composition, curve fitting of the high resolution XPS spectra were acquired in the regions of the C1s, O1s and La3d photoelectron using the SDP v4.1 software®. Structural features of the as-synthesized La_2O_3 films were analyzed by XRD using a Bruker D8 Advanced, Cu K_{α} radiation at 35 kV, 25 mA and a scan rate of 0.021 min^{-1} , whereas their morphology was examined by high-resolution field emission scanning electron microscopy (HR FE-SEM), QUANTA™ 3D FEG. The image was acquired using secondary electrons (SE) at 3.0 kV and at room temperature with a work distance of 6.9 mm ($50,000\times$).

2.2. Optical properties

The optical constants and thickness of the films were obtained by spectroscopic ellipsometry (SE). The measurements were carried out by a Jobin Yvon Uvisel DH10 spectroscopic ellipsometer using an incidence angle 70° and a photon-energy range of 1.5–4.5 eV with an interval step of 0.05 eV. The DeltaPsi2 software® was used to fit the experimental spectra with a three-phases optical model: (1) a crystalline Si layer as the substrate; (2) the metal oxide film modeled as one optical layer (OL) and finally (3) a surface layer modeled as an hybrid La_2O_3 OL with variable void percentage.

UV–vis spectra of the as-prepared samples were recorded in the wavelength range from 200 to 1100 nm using a Perkin Elmer double beam spectrometer, the thickness and refractive index were also calculated using Swanepoel method from the interference fringes region.

The waveguide properties of selected sputtered La_2O_3 films were measured with a prism coupler system (model Metricon 2010, Metricon Corporation, Pennington, USA) which comprises three laser wavelengths of $\lambda = 632.8, 1305$ and 1536 nm. The as-prepared lanthanum films were brought into contact with the base of the TiO_2 prism by means of pneumatically operated coupling head, creating a small gap between the film and the prism. Using a mirror, the laser beam was reflected to the base of the prism where it was totally reflected onto a photodetector. The film/prism arrangement was mounted onto a computer-controlled rotation stage to allow the varied angle of incidence. At certain discrete values of the incident angle θ_m , photons can tunnel across the air gap into the La_2O_3 film propagating as a guided optical mode, which causes a sharp drop in the light intensity reaching the detector.

3. Results and discussion

3.1. Chemical composition, structure and morphology

The chemical composition and depth distribution of the constituent elements in the Lanthanum thin films were obtained by XPS. Fig. 1 a–b shows the typical La 3d doublet. The deconvoluted assignments components of the La 3d levels as well as the calculated relative abundances of lanthanum forming the film are also presented in Table 1. The energy difference between the $3d_{5/2}$ and $3d_{3/2}$ peaks is approximately of 16.8 eV suggesting the formation of lanthanum oxide. In all samples, the La 3d core level spectra split the signals in four peaks at 836.71, 841.17, 853.48 and 857.97, eV that correspond to La $3d_{5/2}$ and $3d_{3/2}$ satellite peaks, respectively [28]. To verify the chemical state of lanthanum in the film the binding energy difference $\text{BE}(\text{O } 1s) - \text{BE}(\text{La } 3d_{5/2})$ was calculated and compared with that reported for reference sample of La_2O_3 [28]. BE difference is commonly employed for chemical bonding characterization due to such difference is a more robust parameter insensitive to BE scale shifts due to dielectric surface charging typical for oxides

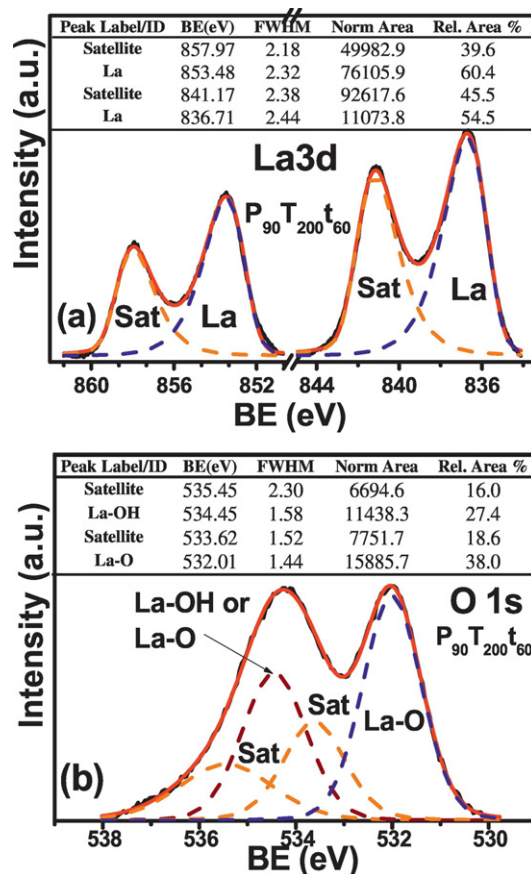


Fig. 1. Fitted high resolution photoemission representative spectra of the a) La 3d and b) O 1s regions using different sputtered conditions, with its corresponding magnitude of the multiple splits.

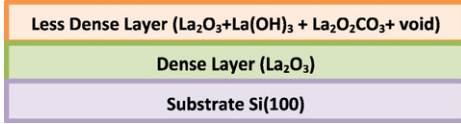
[29–35]. The O 1s envelopes appeared at a glance as being composed of two main peaks located at ~ 532.01 eV accompanied with other at ~ 534.45 eV, the first peak (~ 532 eV) is mainly attributed to La_2O_3 whereas the other peak (~ 534) could be related with hydroxide and a small quantity of carbonate compounds. The O/La ratio (Table 1) and binding energy differences $532-837 \approx 305$ showed that the La–O binding is mainly formed by La_2O_3 , which is in agreement with previous reports [21]. Additionally, it is important to note that lanthanum oxides in hexagonal structure is favored by 0.2 eV/formula unit compare to the cubic, thus we assumed that during deposition process, hexagonal lanthanum initially forms La_2O_3 [36], and then upon exposure to the atmosphere reacts with moisture, leading to extensive surface hydroxylation. It is also clear that sputtered conditions cause a variation in the O/La ratio in the range from 0.67 to 0.77. The chemical environment around the O atoms indicates that, under the synthesis parameters, the films that can be obtained are free of contamination or foreign compounds such as $\text{La}_2\text{O}_2\text{CO}_3$, however, the lanthanum hydroxides or carbonates could be readily formed when the film is exposed to the air [37–38].

Fig. 2 a–b shows the X-ray diffraction patterns of $\text{La}_2\text{O}_3/\text{Si}(100)$ thin films at increasing r. f. power and deposition time. In general, the peak positions and relative intensities in all XRD patterns match well with the combination of hexagonal La_2O_3 (ICDD 05-0602) [37] $\text{La}(\text{OH})_3$ (ICDD 36-1481) [38], and $\text{La}_2\text{O}_2\text{CO}_3$ (ICDD 48-1113) [38]. It is commonly reported for other synthesized methods that lanthanum oxide films are amorphous at low temperatures and eventually crystallized to cubic at 600°C and hexagonal structures after 750°C [25]. However, in our case, the hexagonal phase seems to appear when the films are obtained with a substrate temperature of 80°C and remains at 200°C . The XRD measurements corroborates that lanthanum oxide films were very reactive in the air with formation of lanthanum hydroxides and

Table 1

Composition in atomic percent of elements identified on sputtered Lanthanum thin films, its corresponding O/La ratio, crystallite size and optical parameters calculate from ellipsometry and UV–vis measurements.

Sample	At. % La 3d	At. % O 1s	O/La ratio	ACS La ₂ O ₃ (nm)	ACS La ₂ (OH) ₃ (nm)	ACS La ₂ O ₂ CO ₃ (nm)	Thickness Layer 1 (nm)	Thickness Layer 2 (nm)	Thickness Total (nm)	DR nm s ⁻¹
P ₆₀ T ₂₀₀ t ₂₅	58.8	41.1	0.70	15.2 (1.0)	14.3 (1.1)	15.9 (0.8)	296.7 (9.1)	47.6 (5.5)	344.3	0.229
P ₆₀ T ₂₀₀ t ₄₀	59.3	40.0	0.67	24.2 (5.1)	44.4 (2.5)	50.4 (5.7)	262.9 (6.6)	88.5 (1.2)	351.4	0.146
P ₆₀ T ₂₀₀ t ₆₀	56.9	43.2	0.76	13.5 (4.0)	11.2 (0.2)	14.2 (1.2)	312.4 (6.1)	36.3 (3.7)	348.7	0.097
P ₉₀ T ₂₀₀ t ₂₅	56.9	43.0	0.76	19.3 (1.3)	55.4 (5.4)	69.6 (6.1)	453.1 (9.9)	40.7 (4.6)	493.8	0.329
P ₉₀ T ₂₀₀ t ₄₀	57.4	42.4	0.74	62.3 (3.5)	63.5 (9.6)	43.6 (3.9)	621.7 (10.7)	38.1 (3.3)	659.8	0.275
P ₉₀ T ₂₀₀ t ₆₀	56.4	43.4	0.77	62.4 (2.6)	57.6 (5.6)	54.2 (2.1)	760.4 (19.4)	44.6 (3.8)	805.0	0.224
Sample	Refractive ordinary index (n, at 2 eV)		Refractive extraordinary index (n, at 2 eV)		Ellipsometric measurements			UV–vis measurements		
					Porosity (void) Layer 2	$\omega_{L,E}$	$\omega_{L,O}$	Thickness (nm)	Refractive index (n) at 2 eV	
P ₆₀ T ₂₀₀ t ₂₅	1.744		1.752		17.5 (2.2)	9.2 (0.2)	10.7 (0.4)	–	–	
P ₆₀ T ₂₀₀ t ₄₀	1.761		1.752		11.5 (1.2)	9.2 (0.2)	10.2 (0.2)	–	–	
P ₆₀ T ₂₀₀ t ₆₀	1.738		1.718		15.0 (1.8)	9.9 (0.2)	10.3 (0.4)	301.9	1.558	
P ₉₀ T ₂₀₀ t ₂₅	1.704		1.678		12.4 (1.9)	9.2 (0.02)	10.3 (0.1)	268.9	1.858	
P ₉₀ T ₂₀₀ t ₄₀	1.722		1.706		22 (2.2)	9.6 (0.2)	9.96 (0.1)	273.9	1.776	
P ₉₀ T ₂₀₀ t ₆₀	1.704		1.690		31.3 (3.1)	9.3 (0.2)	10.3 (0.2)	613.7	1.719	



ACS. Average crystallite size. DR, deposition rate.

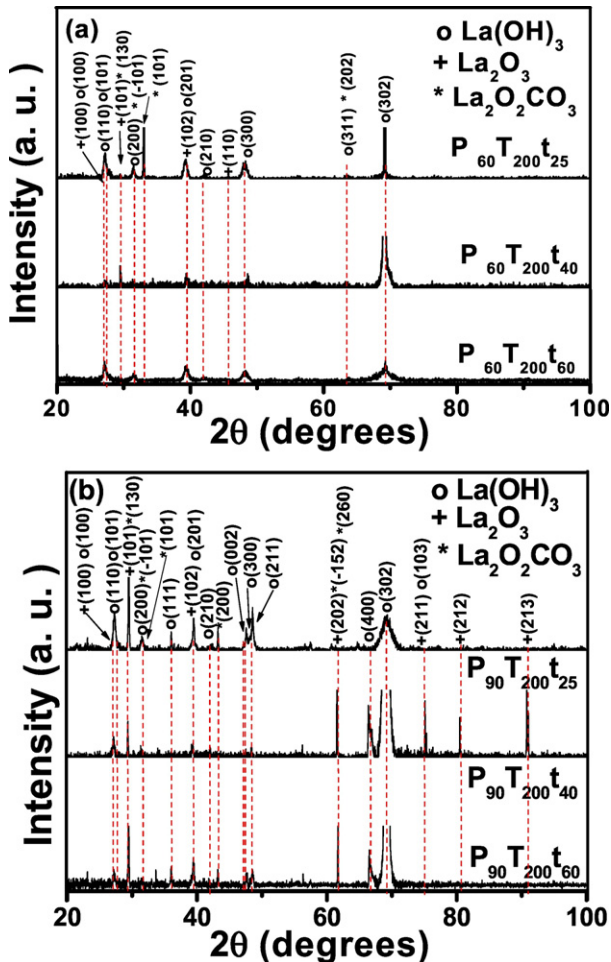


Fig. 2. XRD diffractograms of La₂O₃ films obtained at (a) 60 and (b) 90 W onto Si (100) substrates.

carbonates [26]. Thus, the transformation of h-La₂O₃ to h-La(OH)₃ or m-La₂O₂CO₃ can be obtained when the La₂O₃ coatings are exposed to the air. The different power, time and substrate temperature cause differences in the intensity and broadening of the peaks; which is correlated with the thickness or particle size of the coating. The deviations are particularly observed at (100), (101), (102), (110), (202), (211), (212), (213) planes for h-La₂O₃; (100), (110), (101), (200), (111), (201), (210), (002), (300), (311), (211), (400), (302), (103) for h-La(OH)₃ and (130), (−101), (101), (202), (200), (−152), (260) planes of the monoclinic structure of La₂O₂CO₃. The average crystallite size was estimated by the Scherrer's equation $L = 0.9\lambda/\beta\cos\theta$, where λ (0.154 nm) is the wavelength of X-ray-beam, β is the full-width at half maximum (FWHM) and θ is the Bragg angle (Table 1), using the width of the (100) and (102) planes for La₂O₃; (110), (101), (200), (201), (210), (300) and (302) planes for La(OH)₃. Additionally, (130) and (−101) planes were used to estimate crystallite size for La₂O₂CO₃. It is well known that Scherrer method is a rough estimation due to broadening peaks can be the result of other factors such as microstrains (faulting, dislocations and grain surface relaxation) or even instrumental peak profile among others. Thus, the values showed in Table 1 were obtained from a scheme of deconvolution analysis of each XRD peaks, in order to separate the broadening of each compound, followed by a Gauss fitting. Finally, the crystallite size and uncertainties for the crystallite diameter was calculated using the following equation, which was obtained by propagating the error for λ , β and θ in the Scherrer formula.

$$\frac{\Delta L}{L} = \sqrt{\left(\frac{\Delta\lambda}{\lambda}\right)^2 + \left(\frac{\Delta\beta}{\beta}\right)^2 + (\tan(\theta)\Delta\theta)^2} \quad (1)$$

The error for the X-ray wavelength was obtained from the diffractometer user manual. The error for both β and θ were obtained from the Gaussian fit to the experimental data. The crystallite size was in the range of 11.2 (0.2) to 69.6 (6.1) nm depending on the deposition conditions. It is obvious that the crystal size in all deposited films is dependent on deposition parameters, i.e. the effect of sputtering power and deposition time on the crystal size was not clear. It was expected

that higher power provided energy of the sputtered particles to enhance surface mobility and, consequently, it improves film quality; however, this mobility also causes an alterations of tensile and compressive forces (microstrains; dislocations, vacancies, shear planes) or difference in the incident energy levels resulting in a contracting or broadening of the diffraction peaks.

3.2. Optical properties

Considering an hexagonal structure of the as-prepared La_2O_3 thin films and the rapid transformation in air of $\text{h-La}_2\text{O}_3$ to $\text{h-La}(\text{OH})_3$ or $\text{m-La}_2\text{O}_2\text{CO}_3$, a light beam passed through the crystal, the beam splits into two rays, with two characteristic an anisotropic model was used to determinate ellipsometric parameters. In anisotropic samples, (Ψ, Δ) values vary according to the orientation of the index ellipsoid. In this case, when refractive indices, ordinary (constant) for any direction in the crystal and extraordinary (variable) that depends on the direction. The ellipsometric parameters I_s and I_c measured at $\varphi = 70^\circ$ and its corresponding fitting are shown in Fig. 3a (60 W deposited samples) and Fig. 3b (90 W deposited samples).

However, it should be mentioned that similar results (i.e. the refractive index values and thickness were the same) were also obtained assuming a linear gradient of the refractive index, where the optical layer is modeled as top-bottom refractive indices.

The optical layer was parametrized using a single Lorentz oscillator for dielectric materials model (Eq. (2)), considering anisotropy in the optical properties due to the film composition and the fact that although the films are polycrystalline, they grow with a preferred orientation.

$$\varepsilon = \varepsilon_\infty + \frac{(\varepsilon_s - \varepsilon_\infty)\omega_t^2}{\omega_t^2 - \omega^2 + i\Gamma_0\omega} \quad (2)$$

where ε is the complex dielectric function, ε_∞ and ε_s are the high-frequency and static dielectric constants, ω_t is resonant frequency of the oscillator whose energy corresponds to the absorption peak. Γ_0 is the broadening of the oscillator also known as the damping factor and ω is photon energy.

In Fig. 3 a–b, the solid lines represent the simulated values for photon energies from 1.5 to 4.5 eV, while the symbols represent the measured data. In general, the model fits very well, although small discrepancies between the measured and theoretical values can be observed, which are also reflected in a higher figure of merit, as can be observed in Table 1. The ellipsometric measurements clearly support a two-layer structured for the sputtered films. The best fittings were reached using two layer model i.e. a La_2O_3 dense layer followed by a less dense layer composed of $\text{La}_2\text{O}_3 + \text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3 + \text{voids}$. The thickness of La_2O_3 layer increased with both power and deposition time. However, the dependence of the deposition rate on time at low power (60 W) is very weak. On the other hand, the thickness of the samples deposited at 90 W almost doubled when the deposition time was increased from 25 min to 60 min (Fig. 4). Thick films were obtained at $P_{90}T_{200}t_{60}$ with an interface layer (La_2O_3 dense layer) with a thickness of 760.4 nm combined with a less dense layer ($\text{La}_2\text{O}_3 + \text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3 + \text{voids}$) of 44.6 nm; whereas for the thinnest films at $P_{60}T_{200}t_{40}$, the best fit was achieved with 296.7 nm for the dense layer, followed with a thickness of 47.6 nm for the outer layer ($\text{La}_2\text{O}_3 + \text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3 + \text{voids}$). Thus, the $\text{La}_2\text{O}_3/(\text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3)$ layer thickness fluctuates from ~344.3 to 805.0 nm and it is mainly affected by power and deposition time.

The refractive index dispersion, $n(E)$ of the dense $\text{La}_2\text{O}_3 + \text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3$ layers fabricated at two different powers values are shown in Fig. 5a–b. In the figure, both ordinary and extraordinary refractive indexed were plotted. Refractive indices around 1.66–1.90 were obtained, in good agreement with values reported in the literature [39].

The polarizability of the films plays a key role for the determination of the size-dependent properties of dielectric materials as well as for the

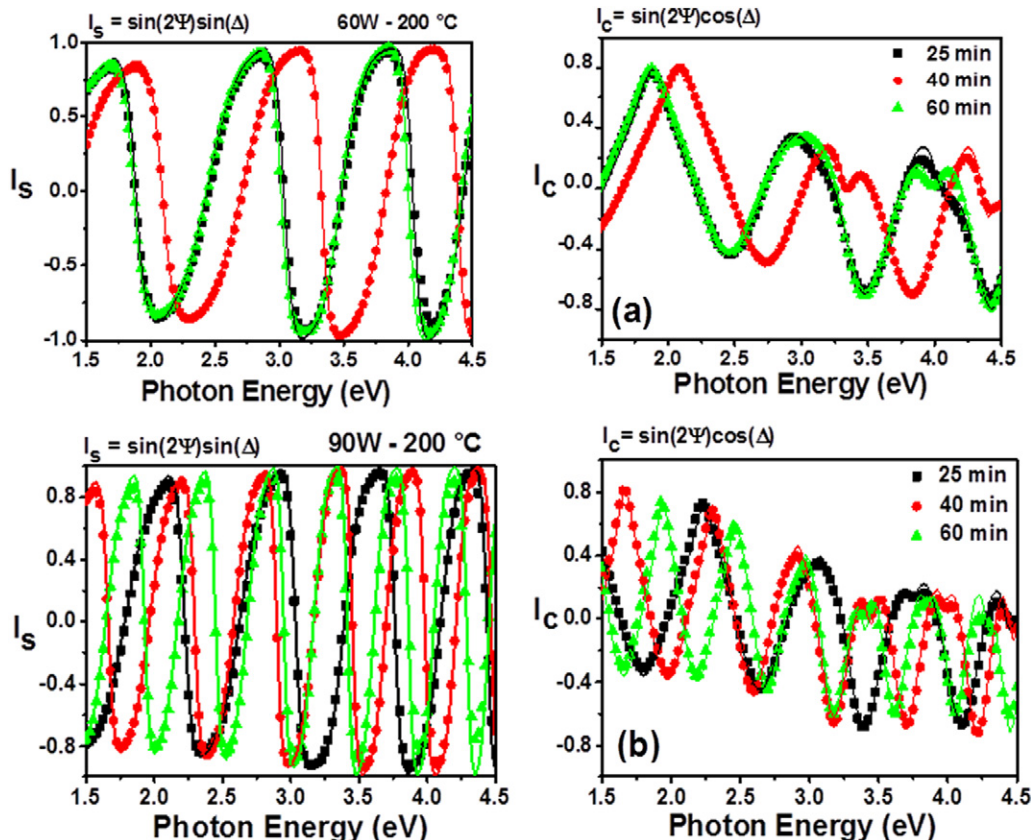


Fig. 3. Optical intensities I_s and I_c showing the best fit for sputtered lanthanum oxide films on silicon substrates.

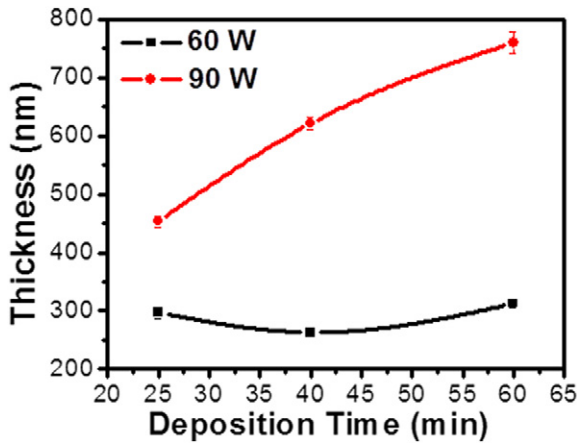


Fig. 4. Dependence of thickness with deposition time of La_2O_3 thin films at 60 and 90 W.

development of new optical components for polarized light with an adequate refractive index. Thus, the variations have been associated with the densification (morphology) and crystallinity of the samples [40].

In order to validate the ellipsometric model and the fitted data, a comparison was realized with UV–vis measurements using as-deposited films on glass substrates (Fig. 6). Two distinct wavelength regions can be observed in the spectra; at longer wavelengths (>350 nm), transmission exhibits oscillation from interference effects in the transparent films; whereas at shorter wavelengths (<350 nm), transmission rapidly decreases to zero. The average transmittance of the samples was over 76% in the visible wavelength region. In general, it is noticed for the samples that the average transmittance in the near infrared (NIR) spectral range tends to achieve the glass transmittance value (90%). The measured transmittance spectra for samples that displayed interference fringes were used to calculate the optical constants applying the method

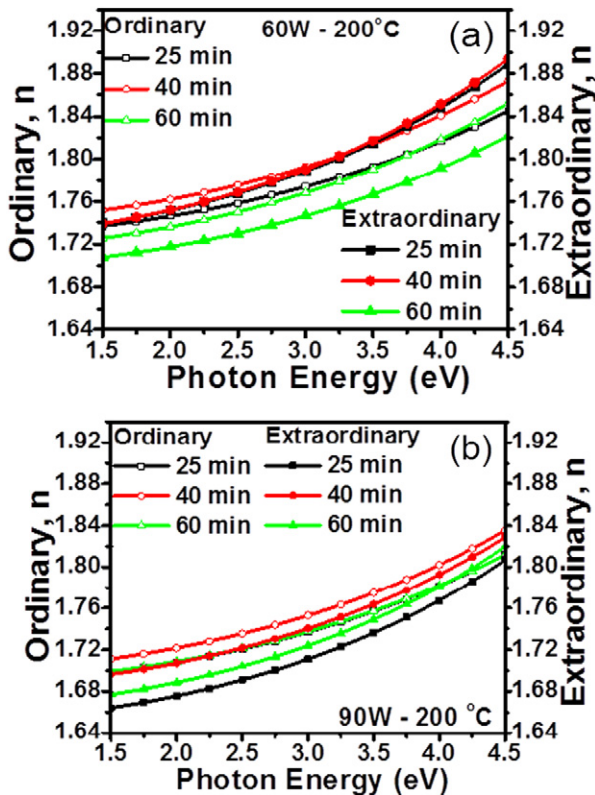


Fig. 5. Ordinary and extraordinary refractive indexes profiles determined by proposed model as function of photon energy.

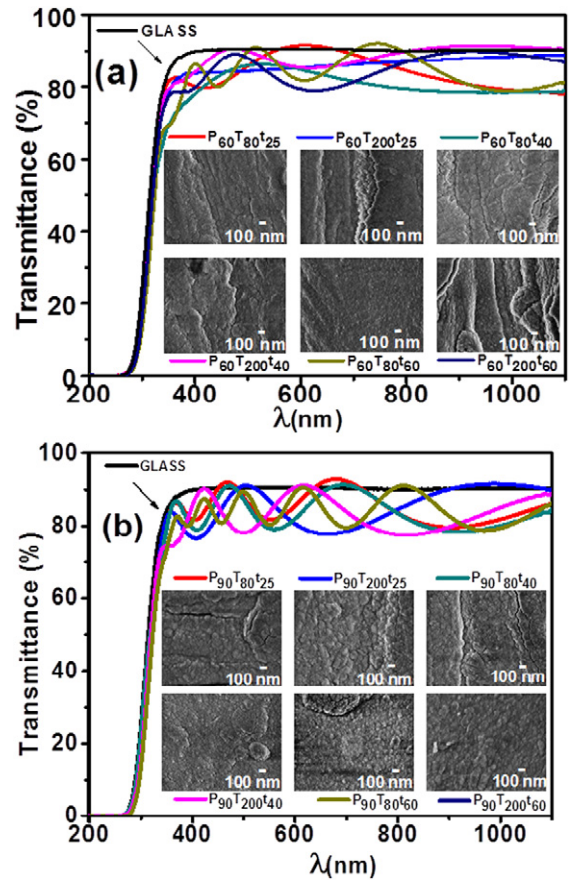


Fig. 6. UV-vis spectra of the as-prepared samples a) 60 W and b) 90 W.

suggested by Swanepoel [41]. The approach is based on the use of maximal and minimal interference effects in the transparent region of the spectrum for calculation of the refractive index and film thickness [42–45]. The calculated optical properties for UV–vis measurements are shown in Table 1. Accurate values of n and thickness were only considered during calculations whereas other solutions were considered out of the range. Due to the frequency dependence of these parameters at the infrared region, the data were calculated at 2 eV. The results exhibit good agreement between ellipsometric and UV–vis measurements. The refractive index at 2 eV ranged from 1.558 to 1.858 demonstrating a strong dependence on the deposition parameters with thickness varied from 268.9 to 613.7 nm. SEM inset micrographs observed at $50,000\times$, show a densely packed surface, with small pores, and smooth domes that covers uniformly the surface substrate. Some experimental conditions presented small cracks in the films that tend to disappear with the power and deposition time. The growth process is in agreement with previous works that indicate that the kinetic energy of the expelled particle increases with the power, when the atoms fall on the surface of the substrate are adsorbed and the film is densified. Under these conditions, the lanthanum nanoparticles coalesce and the size gradually increases until a continuous film is formed and the surface becomes irregular. The evolution of the lanthanum film can be explained in terms of the Stranski-Krastanov model, where the growth changes from layers to islands after a monolayer or two due to the change in energy with the successive layer. The small differences in the morphology may explain the discrepancies in the optical parameters.

Fig. 7 shows the TE guided modes measured at different wavelengths (633, 1305 and 1536 nm), which propagate in the $\text{La}_2\text{O}_3 + \text{La}(\text{OH})_3 + \text{La}_2\text{O}_2\text{CO}_3$ film synthesized at $\text{P}_{90}\text{T}_{200}\text{t}_{60}$. The sharpness of the intensity drops indicates a good confinement of the light into the film and also the smoothness of the film surface. At shorter wavelength, two TE propagation modes can be clearly identified, enabling

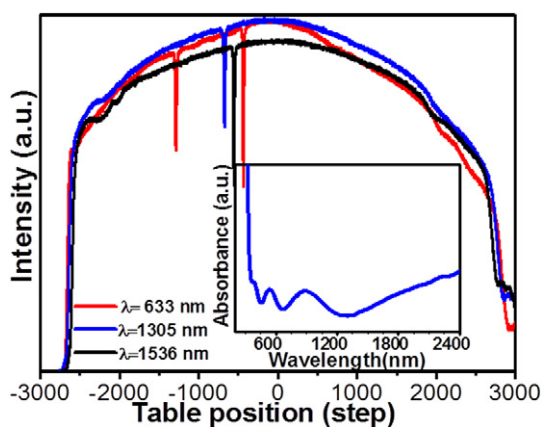


Fig. 7. Propagation modes for La_2O_3 films obtained at $P_{60}\text{T}_{200t_{60}}$ and measured with the prism coupling technique at three different wavelengths 633 nm; 1305 nm and 1536 nm and its comparison with UV-vis spectra.

the simultaneous calculation of refractive index and film thickness. For longer wavelengths only one mode propagates along the film. Assuming the thickness calculated at wavelength of 633 nm (636.1 nm), it was possible to obtain the refractive index at 1305 nm (1.957) and 1536 nm (2.011). The values reveal a refractive index that increase as wavelength increases from 1305 to 1536 nm; thus the negative refractive index dispersion (anomalous dispersion) is a consequence of an adsorption band that the material exhibits in this region (see inset in Fig. 7). This behavior is in agreement with the decrease in the film transmittance observed in spectra acquired in the NIR region 1000 to 1500 nm. It is important to highlight that for comparison calculate thickness film by prism coupling (636.1 nm) and Swanepoel method (613.7 nm) are quite similar, however are far from the final average thickness determine by ellipsometry (805 nm).

Samples obtained at lower power ($P_{60}\text{T}_{200t_{60}}$) were also characterized with this optical method. It is worth notice that under these conditions thickness of the film obtained by prism coupling (342 nm) and Swanepoel method (302 nm) well matched with ellipsometric measurements (348 nm); however, there is a discrepancy in the refractive indexes (1.975 at 2.0 eV and 1.558 at 2.0 eV), whereas in the case of ellipsometry the ordinary and extraordinary indexes are about (1.738–1.718 at 2.0 eV). Such discrepancies might be related to inhomogeneity of the film thickness, which became critical when techniques with different spatial resolution are compared.

4. Conclusions

In this work, we report the influence of the synthesis parameters on the optical properties of sputtered La_2O_3 thin films. The XRD results revealed that the crystallite size was in the range of 10–65 nm. The lanthanum films were deposited as La_2O_3 hexagonal structure and when films are exposed to air, the oxide surface film is transform into $\text{h-La}(\text{OH})_3$. The XPS analyses corroborate the transformation of $\text{h-La}_2\text{O}_3$ to $\text{h-La}(\text{OH})_3$ and stoichiometry of the compounds. The application of r. f. magnetron sputtering as a preparation method of La_2O_3 thin films can be used to modulate the optical properties from tailored synthesis parameters. The high refractive index ($n \sim 2$) and high transmittance in the visible region provide the interesting properties for the lanthanum oxide films as candidate to develop new optical active devices.

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