

# On the Structure Factors of Aqueous Mixtures of 1-Propanol and 2-Propanol: X-Ray Diffraction Experiments and Molecular Dynamics Simulations

José G. Méndez-Bermúdez,\* Hector Dominguez, László Temleitner,  
and László Pusztai

The structure factor of pure 1-propanol, 2-propanol, and mixtures of 1-propanol/water and 2-propanol/water, as a function of composition, has been determined experimentally and by molecular dynamics simulations. The primary aim is to find interatomic potentials that reproduce measured structural data at the highest possible level. For this reason, various alcohol potential models have been employed, including united atom (UA) and all atom (AA) types, in combination with a TIP4P-based model for water. In order to improve agreement with experimental values of the dielectric constant and mass density, a new UA force field for the alcohols has also been constructed. In terms of structural properties, the AA model reproduces experimental results better than any of the UA models for all compositions.

## 1. Introduction


1-Propanol (*n*-propanol, or simply “propanol,”  $\text{CH}_3\text{-CH}_2\text{-CH}_2\text{-OH}$ ) and 2-propanol (isopropanol,  $\text{CH}_3\text{-CH}(\text{-OH})\text{-CH}_3$ ) are among the few mono-alcohols that are fully miscible with water under ambient conditions. They also have their industrial uses, at the million tons/year extent, mostly as solvents.

Dr. J. G. Méndez-Bermúdez  
Centro Universitario de los Valles  
Universidad de Guadalajara  
Carretera Guadalajara-Ameca Km 45.5, 46600 Jalisco, México  
E-mail: jose.mendez@valles.udg.mx

Dr. H. Dominguez  
Instituto de Investigaciones en Materiales  
Universidad Nacional Autónoma de México  
Circuito Exterior, 04510 Ciudad de México, México

Dr. L. Temleitner, Dr. L. Pusztai  
Wigner Research Centre for Physics  
Hungarian Academy of Sciences  
Konkoly Thege út 29-33, H-1121 Budapest, Hungary

Dr. L. Pusztai  
International Research Organization for Advanced Science and  
Technology (IROAST)  
Kumamoto University  
2-39-1, Kurokami, Chuo-ku, Kumamoto 860-855, Japan

 The ORCID identification number(s) for the author(s) of this article can be found under <https://doi.org/10.1002/pssb.201800215>.

DOI: 10.1002/pssb.201800215

Concerning the microscopic structure of the pure substances, a host of diffraction studies have been conducted over the past decades. Pure 1-propanol has been considered by X-ray<sup>[1–4]</sup> and neutron diffraction,<sup>[5–8]</sup> whereas neutron diffraction has been the preferred method of investigation for pure 2-propanol.<sup>[8–10]</sup> For aqueous mixtures of 1- and 2-propanol, only a few X-ray<sup>[11–13]</sup> and neutron<sup>[14]</sup> diffraction measurements have been performed; as an interesting example, the work of Tong et al.<sup>[15]</sup> is worth mentioning, which applied infrared spectroscopy for observing conformational changes in mixtures of 1-propanol and water.

According to the best of our knowledge, only very few computer investigations compared simulation results to diffraction data in the reciprocal space: such comparisons are available for liquid pure 1-propanol only.<sup>[3,4,16]</sup> On the other hand, in a recent study,<sup>[17]</sup> molecular dynamics computer simulation results have been contrasted to experimental data on the temperature-dependent dielectric constant of 1-propanol/water mixtures. In most cases, agreement between simulation and experiment was excellent for one specific type of potential (see ref. [17] for details). It would be important to learn whether the same (or similar) interatomic potentials were able to describe a most basic property, the liquid structure, as well. The primary focus of the present work then is to establish if an interatomic potential model that reproduces the measured dielectric constant of 1-propanol/water mixtures can, at the same time, provide structural models that are consistent with the experimental X-ray structure factors.

In addition, there still is a controversy concerning even the neutron diffraction total structure factor of 2-propanol, as it is exemplified by ref. [8] (Figure 5 of that publication), particularly at the lower scattering variable,  $Q$ , range (below ca.  $3 \text{ \AA}^{-1}$ ). This is also why a comparative computer simulation study has been missing: we are now providing simulated X-ray and neutron-weighted total structure factors, along with simulated densities of pure 1- and 2-propanol and their mixtures with water.

## 2. Experimental Section

Synchrotron X-ray diffraction measurements have been carried out at the BL04B2<sup>[18]</sup> high energy X-ray diffraction beamline of

**Table 1.** Original<sup>[22]</sup> and modified (UA\_new) UA potential values of partial charges and Lennard–Jones (LJ) parameters for 2-propanol.

2-Propanol sites	UA charges (e) <sup>[22]</sup>	UA_new charge (e)	UA energy $\epsilon$ LJ parameter (kJ/mol) <sup>[22]</sup>	UA_new energy $\epsilon$ LJ parameter (kJ/mol)	UA $\sigma$ LJ parameter (nm) <sup>[22]</sup>	UA_new $\sigma$ LJ parameter (nm)
H(-OH)	0.4350	0.4667	0.00000	0.00000	0.000	0.00000
O(-OH)	−0.700	−1.2088	0.77325	0.55674	0.302	0.31861
C(C-OH)	0.2650	0.8485	0.08314	0.05986	0.433	0.45682
C(CH <sub>3</sub> -C)	0.0000	−0.0532	0.81482	0.58667	0.375	0.39563

the Japan Synchrotron Radiation Research Institute (SPring-8, Hyogo, Japan). The incoming wavelength was 0.2023 Å (corresponding to a photon energy of 61.3 keV). Diffraction patterns from the samples were recorded in the transmission mode by a single HPGe detector between scattering variable,  $Q$ , values of about 0.2 and 16 Å<sup>−1</sup>. Samples were put in 2 mm diameter, thin-walled quartz capillaries (GLAS Müller, Germany). The capillaries were mounted in the automatic sample changer. Diffraction patterns were recorded in three overlapping frames that differed by the width of incoming beam.

Measured raw intensities were normalized by the incoming beam monitor counts, then corrected for absorption, polarization, and contributions from the empty sample holder. Finally, diffraction patterns over the entire  $Q$ -range were obtained by normalizing and merging each frame in electron units, then removing inelastic (Compton) scattering contributions following a standard procedure.<sup>[19]</sup>

Since each sample was considered also for neutron diffraction experiments, heavy water, D<sub>2</sub>O, and heavy alcohols, C<sub>3</sub>D<sub>8</sub>O, have been used for preparing the mixtures. Due to the need for costly deuterated compounds, we have not been able to cover the entire concentration range for both 1- and 2-propanols. In the end, the two pure alcohols, as well as mixtures with 10 and 50 mol.% of (fully deuterated) 1-propanol, and 20 and 80 mol.% of (fully deuterated) 2-propanol could be prepared. We hope that it is correct and fair to note that both the  $Q$ -range and the statistical accuracy of the new data surpass those of the previously determined X-ray total structure factors – this is a justification why new measurements were timely. (Neutron diffraction measurements and data analyses are still planned for the future.)

### 3. Molecular Dynamics Simulations

Molecular dynamics computer simulations have been performed by the GROMACS<sup>[20]</sup> software (version 4.5). Water/alcohol mixtures (1-propanol/water and 2-propanol/water) were

constructed at different compositions using two alcohol potential models, combined with the TIP4P/ε water potential<sup>[21]</sup>; this water model was originally optimized for providing the correct dielectric constant of pure water. The first alcohol model considered all atoms (AA), including each hydrogen, explicitly. In the other potential model, CH<sub>n</sub> (n: 1, 2, 3) groups were taken as single sites: this is the philosophy of the united atom (UA) approach. For 1-propanol we took the UA force field parameters already reported in our previous work.<sup>[17]</sup>

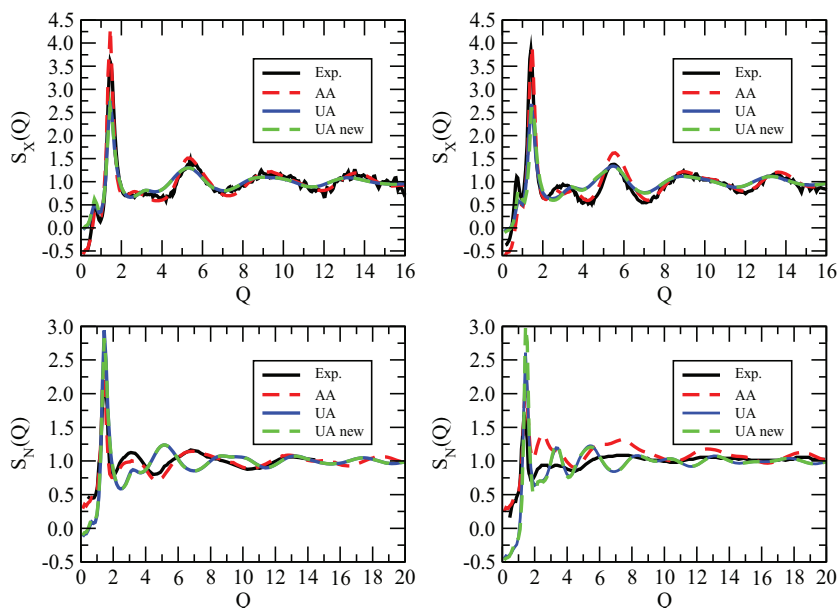
Simulations containing 2-propanol were carried out using parameters reported in the literature for the UA<sup>[22]</sup> and AA<sup>[23]</sup> types. In addition, UA parameters from the literature have been re-parametrized, similarly to the 1-propanol case.<sup>[17]</sup> That is, the new 2-propanol UA parameters were adjusted to obtain the correct experimental dielectric constant of pure 2-propanol. The new Lennard–Jones (LJ) parameters were calculated by conducting several simulations, trying several LJ parameter pairs, to fit the experimental surface tension and density at temperature 298 K and pressure 1 bar. Details of the re-parametrization procedure are given in our previous work<sup>[17]</sup> and in ref. [24]. The original and re-parametrized (marked in short as “UA\_new”) potential parameters are provided in **Table 1**.

In **Table 2**, selected properties of pure 2-propanol for the original and new UA potentials are compared with the corresponding experimental data. The static dielectric constant and the density values are dramatically improved after the application of the re-parametrized UA model. On the other hand, the surface tension slightly departs from the experimental value: this is an indication already that no single interaction potential model is able to re-produce every property of the liquids considered here at the highest level.

All MD simulations were conducted in the NPT ensemble, using the Berendsen barostat and Berendsen thermostat,<sup>[25]</sup> to keep pressure and temperature constant during the entire calculation. Relaxation time constants were fixed at  $\tau_T = 0.1$  ps and  $\tau_P = 2.0$  ps for the temperature and pressure, respectively. Periodic boundary conditions were imposed in all directions.

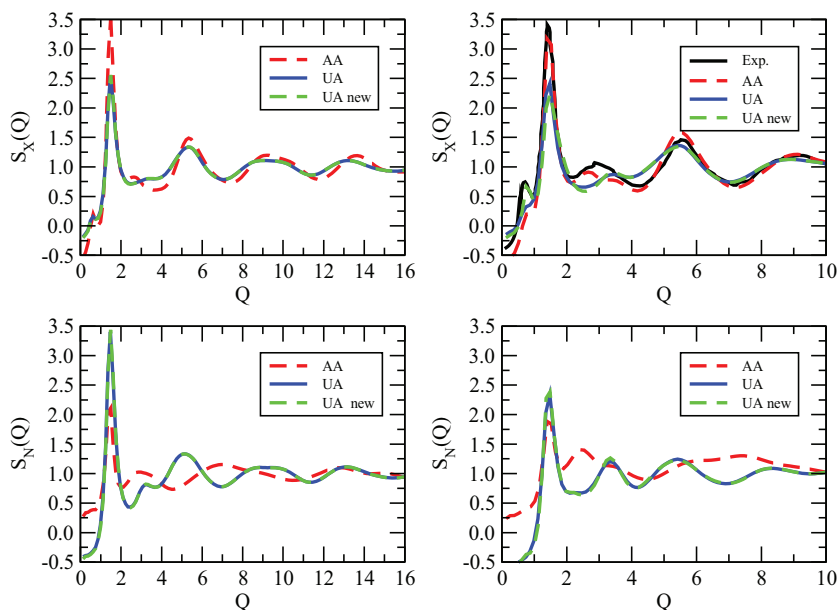
**Table 2.** Various properties of 2-propanol at  $T = 298$  K as calculated via the original<sup>[22]</sup> and re-parametrized (UA\_new) UA force fields. For comparison, the corresponding experimental values are also shown.

Optimized parameters	UA (% error)	UA_new (% error)	Experimental values
Dielectric constant, $\epsilon$	9.97954 (48.17)	19.1433 (0.58)	19.255 <sup>[26]</sup>
Surface tension, $\gamma$ (mN m <sup>−2</sup> )	21.5746 (1.67)	23.6297 (11.36)	21.22 <sup>[27]</sup>
Density, $\rho$ (kg m <sup>−3</sup> )	787.866 (0.86)	781.519 (0.05)	781.11 <sup>[28]</sup>



**Figure 1.** X-ray and neutron-weighted total structure factors,  $S_X(Q)$  and  $S_N(Q)$ , respectively, for pure 1- (left panels) and 2- (right panels) propanol. The experimental neutron diffraction spectra are taken from refs. [6,9]. The experimental temperature for 2-propanol  $S_N(Q)$  is 250 K, whereas the other structure factors were all measured (and simulated) at 298 K. (Note that the experimental  $S_N(Q)$  for 2-propanol is representative only as far as the positions (but not the intensities) of the extrema are concerned.)

The time-step was 2 fs. LJ interactions were neglected beyond a cut-off radius of 1.5 nm. Pure (100%) alcohol models contained 1500 alcohol molecules, while models with 80, 50, 20, and 10%



**Figure 2.** X-ray and neutron-weighted total structure factors,  $S_X(Q)$  and  $S_N(Q)$ , respectively, for aqueous mixtures of 1- (left panels) and 2- (right panels) propanol, at 80% molar concentration of alcohol. The experimental spectrum of the 2-propanol/water mixtures was determined in the course of the present study.

alcohol contained 1024 alcohol and 280 water, 1000 alcohol and 1000 water, 325 alcohol and 1300 water, and 200 alcohol and 1800 water molecules, respectively. The static dielectric constant was calculated over 200 ns of production runs, whereas other properties, such as densities and diffusion coefficients, were analyzed for the last 5 ns of 50 ns simulations.

## 4. Results and Discussion

From the point of view of the present work, the most important simulated quantity to be compared with its experimental counterpart is the total structure factor,  $S(Q)$ . For a most direct assessment of the interaction potential models considered, the comparison is made in the reciprocal (or “Q-”) space.

In order to obtain the X-ray-weighted total structure factor,  $S_X(Q)$ , one needs to combine the Fourier transforms,  $\widehat{g}_{ij}(Q)$ , of all the partial radial distribution functions,  $g_{ij}(r)$ , according to their contributions to the total. This way we have

$$S_X(Q) = \frac{\sum_{i=1}^n c_i^2 f_i(Q) f_i(Q) \widehat{g}_{ii}(Q) + \sum_{i \neq j}^n c_i c_j f_i(Q) f_j(Q) \widehat{g}_{ij}(Q)}{(\sum_{i=1}^n f_i(Q))^2}$$

where  $n$  is the number species,  $c_i$  are molar fractions of specific sites, and the  $f_i(Q)$  X-ray scattering form factors can be obtained as

$$f_i(Q) = \sum_{i=1}^5 a_i e^{-b_i Q^2} + c$$

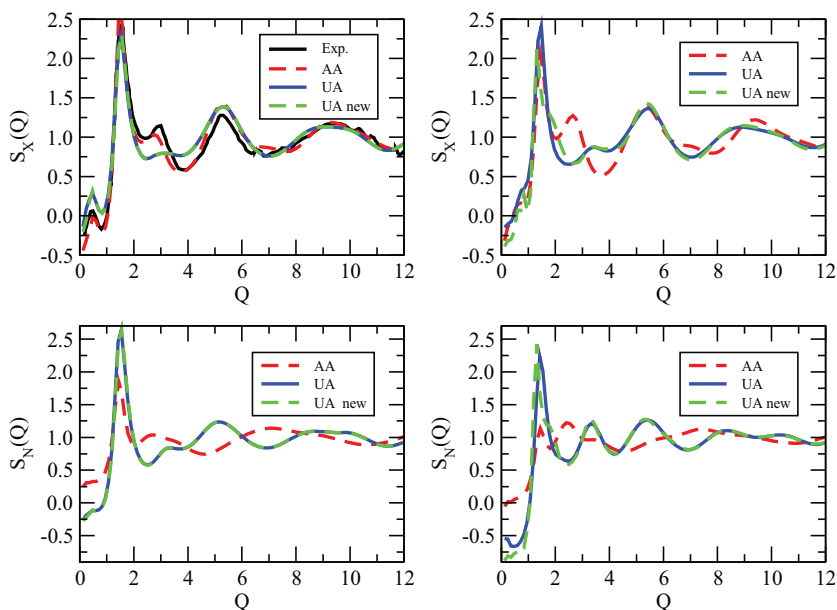
where the  $a_i$ ,  $b_i$ ,  $c$  parameters are taken from ref. [29] for each atom type.

The neutron-weighted total structure factor,  $S_N(Q)$ , is given by the following expression

$$S_N(Q) = \frac{\sum_{i=1}^n c_i^2 b_i \widehat{g}_{ii}(Q) + \sum_{i \neq j}^n c_i c_j b_i b_j \widehat{g}_{ij}(Q)}{(\sum_{i=1}^n c_i b_i)^2}$$

where  $b_i$  are the coherent neutron scattering lengths for each atom type (i.e., AA site).<sup>[30]</sup> (Scattering amplitudes for the UA sites have been approximated by simply summing the corresponding AA scattering lengths.)

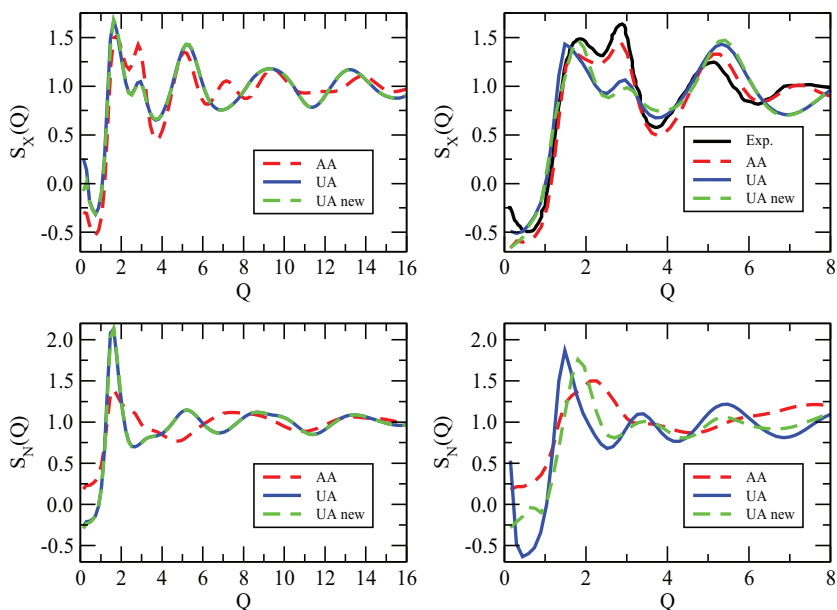
X-ray- and neutron-weighted total structure factors for the 1-propanol and 2-propanol mixtures are displayed in Figure 1–5. Calculated total structure factors are shown for the original<sup>[22]</sup> and re-parametrized (ref. [17] for 1-propanol) UA, as well as for the AA<sup>[23]</sup> alcohol potential models. X-ray experimental



**Figure 3.** X-ray and neutron-weighted total structure factors,  $S_X(Q)$  and  $S_N(Q)$ , respectively, for aqueous mixtures of 1- (left panels) and 2- (right panels) propanol, at 50% molar concentration of alcohol. The experimental spectrum of the 1-propanol/water mixtures was determined in the course of the present study.

data, where available, are from the present work, whereas neutron diffraction data for pure 1- and 2-propanol were taken from refs. [6,9], respectively. (Note that in **Figure 1**, the neutron-weighted experimental structure factor,<sup>[9]</sup>  $S_N(Q)$ , looks somewhat off-scaled. We have attempted to find a proper scaling

observed for 2-propanol (the AA model seems to be slightly better than the two UA models). The good performance of the UA models means that for pure 1- and 2-propanol, we do have structural models (without non-alcoholic hydrogen atoms though) whose dielectric constants and X-ray-weighted total structure factors are nearly equally realistic.



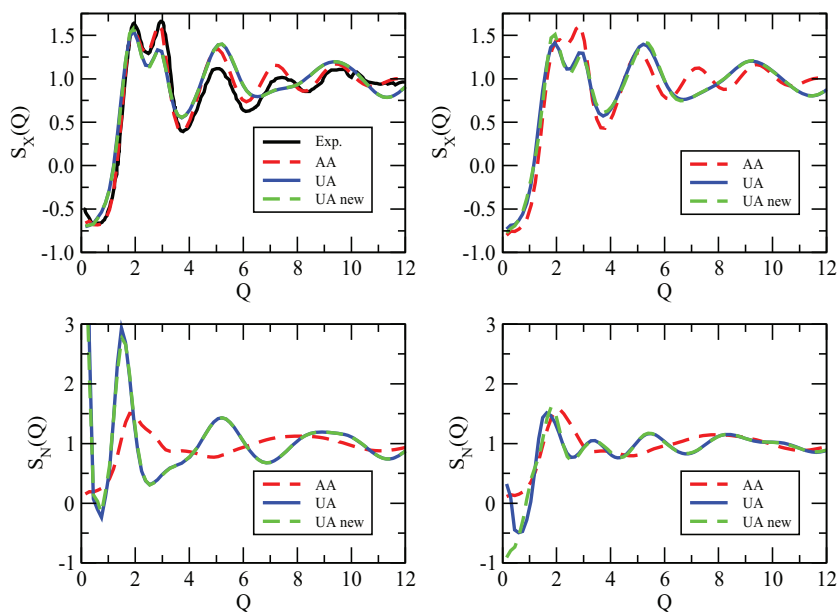
**Figure 4.** X-ray and neutron-weighted total structure factors,  $S_X(Q)$  and  $S_N(Q)$ , respectively, for aqueous mixtures of 1- (left panels) and 2- (right panels) propanol, at 20% molar concentration of alcohol. The experimental spectrum of the 2-propanol/water mixtures was determined in the course of the present study.

factor, but as this effort proved unsuccessful, the original scale has been maintained. It is still obvious that it is the AA model that follows maximum and minimum positions nearly perfectly, where the two UA potentials led to observable shifts in terms of the positions of the extrema.)

As a general statement, for all compositions the AA alcohol model provided a better description of the experimentally determined functions than the other two models, UA<sup>[22]</sup> and UA\_new. Another general observation is that UA and UA\_new have brought about almost identical behavior in terms of the structure (this, again, valid across the entire concentration region). That is, despite the enormous differences in terms of, e.g., the dielectric constants calculated from the two UA type potentials (see Table 2), re-parametrization has not helped in producing better structural models.

Figure 1 shows total structure factors obtained by X-ray and neutron diffraction for 1-propanol (left panels). The simulated  $S_X(Q)$  for all (UA, UA\_new, and AA) alcohol potential models seem to fit the experimental spectrum quite well. Similar behavior is observed for 2-propanol (the AA model seems to be slightly better than the two UA models). The good performance of the UA models means that for pure 1- and 2-propanol, we do have structural models (without non-alcoholic hydrogen atoms though) whose dielectric constants and X-ray-weighted total structure factors are nearly equally realistic.

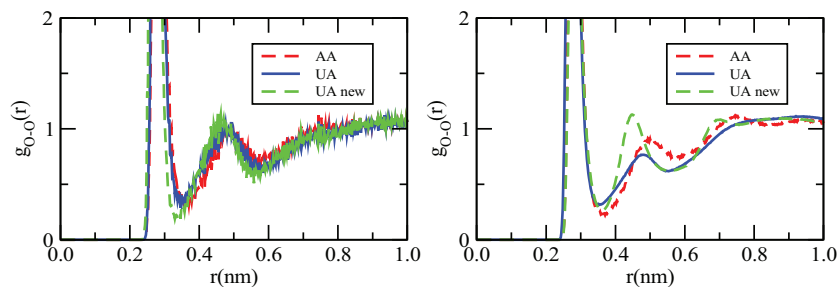
The simulated  $S_N(Q)$  spectrum for AA 1-propanol also provides a good description of the neutron diffraction measurement.<sup>[6]</sup> For 2-propanol, again, the AA model is clearly better than the UA models (cf. also the notes concerning the corresponding experimental data<sup>[9]</sup> above). The misfit between the UA models and experiment is understandable by noting that whereas X-rays are relatively insensitive to hydrogen atoms, the scattering power of (particularly heavy) hydrogen is considerable for neutrons. By uniting the  $CH_n$  groups according to the UA philosophy, the measured information from pair correlations involving H is lost (or at least, severely averaged out); this is why UA-type models are not good choices for describing neutron diffraction experimental data in general. Having this in mind, it is not surprising that in each of Figure 1–5, differences between predictions from UA and AA potentials are large in terms of  $S_N(Q)$ . Also, it is always the AA predictions that may be considered as better in this sense, c.f., for instance, the complete failure of UA models at low  $Q$  values



**Figure 5.** X-ray and neutron-weighted total structure factors,  $S_X(Q)$  and  $S_N(Q)$ , respectively, for aqueous mixtures of 1- (left panels) and 2- (right panels) propanol, at 10% molar concentration of alcohol. The experimental spectrum of the 1-propanol/water mixtures was determined in the course of the present study.

at low alcohol concentrations (see **Figure 4** and **5**). In what follows, more detailed discussions will concern X-ray data only (also for the reason that they are the ones that have been determined during the present study).

The UA models predict  $S_X(Q)$  better at the lowest alcohol concentrations and for pure alcohols than at intermediate (20–80 mol.% of alcohol) compositions. A similar observation may be made concerning the performance of the AA potential model, and could be made earlier<sup>[17]</sup> concerning the temperature-dependent dielectric constant of 1-propanol/water mixtures. This is most probably the unfortunate consequence of the fact that, in general, potential models are developed for pure substances (like UA\_new for 2-propanol in the present work). Consequently, if the performance of potential functions developed for pure alcohols provide unsatisfactory results when applied to aqueous mixtures then a separate potential development procedure would be necessary for every mixture. Such an effort, however, is most certainly too expensive and/or impractical.



**Figure 6.**  $g_{O-O}(r)$  (O: alcohol oxygen) in pure 1-propanol (left) and 2-propanol (right).

Turning now to real space structural properties, selected partial radial distribution functions (PRDF),  $g_{ij}(r)$ , are going to be introduced. We do not wish, at this early stage, to scrutinize every single PRDF until a definite conclusion concerning the most appropriate interatomic potentials is reached. We thus report various O–O PRDF-s that a) have reasonably high weighting factors in the X-ray total structure factor (contrary to PRDF-s that contain hydrogen) and b) are related to hydrogen bonding in the neat substances and in the mixtures alike.

**Figure 6** shows the (only) oxygen–oxygen PRDF present in the pure alcohols. The most striking finding is that while for 1-propanol, the curves that belong to the original<sup>[22]</sup> and modified UA, and the AA<sup>[23]</sup> potentials run together, for 2-propanol there are three distinct lines. The largest departure is registered for the modified UA potentials in both propanols. If we now go back to the X-ray-weighted total structure factors (**Figure 1**), an analogous phenomenon is apparent: whereas for 1-propanol, the three potentials bring about nearly identical  $S_X(Q)$ -s, for 2-propanol, it is only the OPLS/AA<sup>[23]</sup> model that

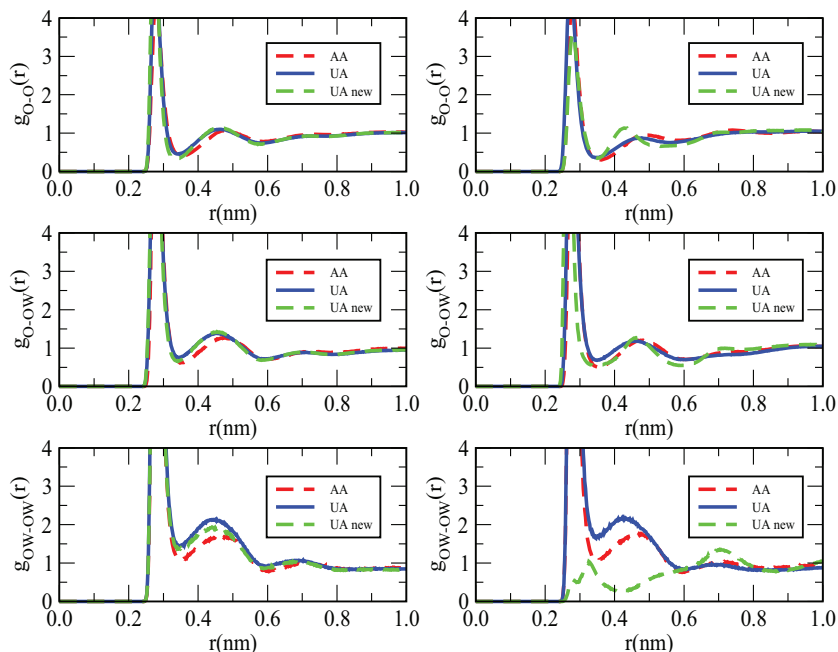
produces a very good agreement with experiment. The consequence of these observations is that while the O–O PRDF-s of 1-propanol originating to all the three potentials applied here can be taken as roughly equally reliable, for 2-propanol only the OPLS/AA force field was able to produce a trustable O–O PRDF.

As it may be discerned from **Figure 2** to **5**, differences between total  $S_X(Q)$ -s originating to the UA-type and AA potential models are somewhat larger for 2-propanol/water mixtures. This is reflected by the corresponding O–O PRDF-s, too, as exemplified by **Figure 7** and **8**, where the three possible O–O partials (alcohol–alcohol, alcohol–water, and water–water) are provided for one alcohol-rich and one water-rich mixture (with 80 and 10 molar % of alcohol). For the 1-propanol containing mixtures, differences observed here-and-there are only in terms of the amplitudes, whereas for 2-propanol/water mixtures variations are more drastic, particularly by the curves corresponding to the modified UA potentials (UA\_new).

In both mixtures, and on each O–O PRDF, a nearest neighbor distance of about 0.3 nm can be observed. This is consistent with the notion that there are alcohol–alcohol, alcohol–water, and water–water hydrogen bonds in these liquids, just as expected.

Another notable feature is that, especially in the case of the alcohol-rich mixture (**Figure 7**), the two pure O–O PRDF-s, alcohol–alcohol and water–water, the functions do not oscillate around unity but display slopes. A similar behavior has been presented (and left un-commented) by a recent, neutron diffraction-based study for 2-propanol/water liquid mixtures.<sup>[14]</sup>





**Figure 7.** Oxygen–oxygen partial radial distribution functions  $g_{O-O}(r)$  (O: alcohol oxygen),  $g_{O-OW}(r)$  (OW: water oxygen), and  $g_{OW-OW}(r)$  for aqueous mixtures with 1-propanol (left) and 2-propanol (right) at 80% molar concentration of the alcohol.

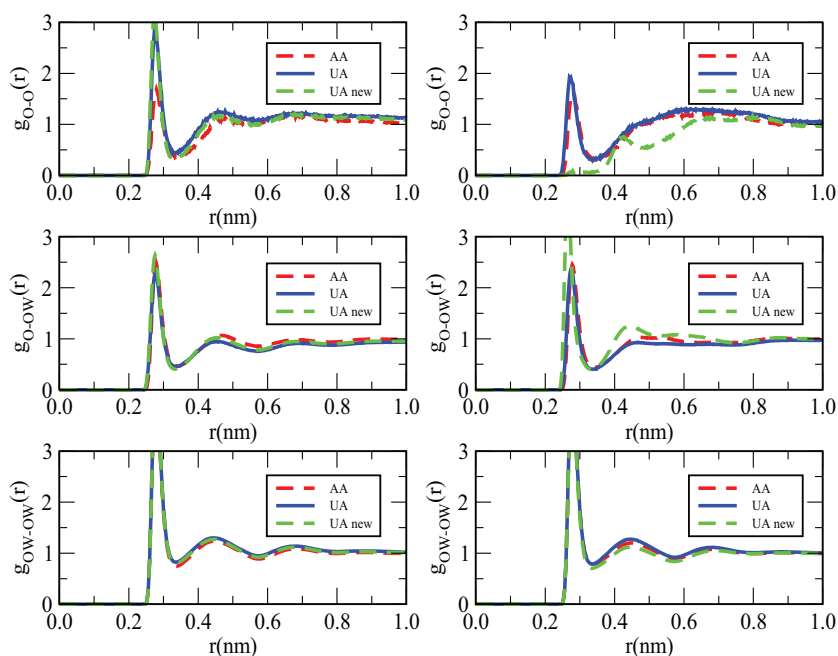
The obvious interpretation would be that water molecules distinctly prefer water neighbors; on the other hand, propanol molecules do not want to have propanol neighbors (the water–water preference is the large effect, cf. Figure 7). Given the

1-propanol mixture. However, at low 2-propanol concentration some differences are depicted, in particular for the UA new model (right top and middle plots of Figure 8). For that model the alcohol oxygen–oxygen interaction is weak, as suggested by the

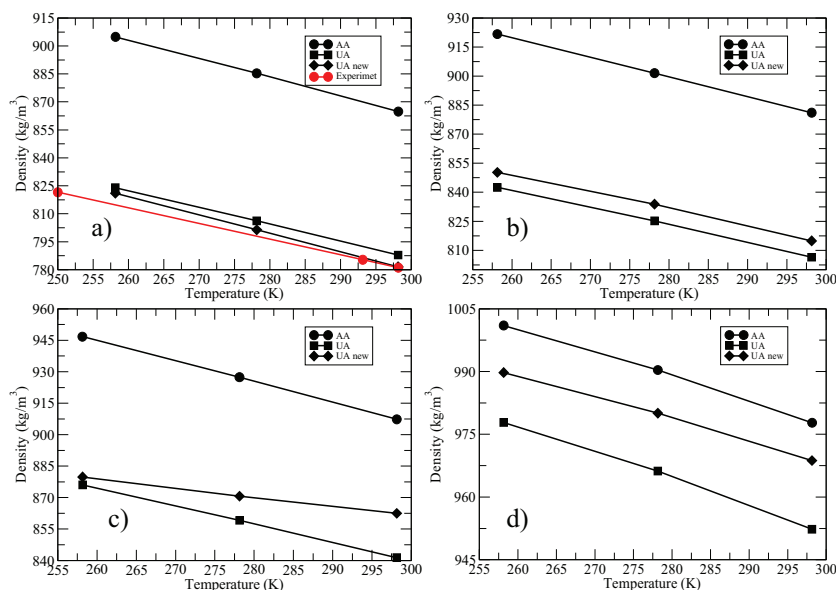
composition (80% of alcohol), this would mean the formation of “water-islands” in the mixture. Once the question of appropriate potential functions is settled, this feature would deserve a separate, in-depth investigation.

From Figure 7 a high first peak for oxygens of 1-propanol is observed, suggesting strong interactions between alcohol molecules at high alcohol concentration. When the alcohol concentration decreases, there are more water molecules around 1-propanol molecules so their head-to-head interaction is reduced, as indicated by a lower first peak in the radial distribution function (see Figure 8). In fact, for the AA model the peak is even lower than for the other alcohol models. It is also interesting to observe a higher first peak in the  $g(r)$  of water (oxygen) with oxygen of the 1-propanol at high alcohol concentration than seen at low concentration (middle plots of Figure 7 and 8). At high alcohol concentration, water and (alcohol) OH groups display stronger attraction, whereas at low alcohol concentration there are fewer water molecules close to the (alcohol) OH groups. Similar trends are observed for the water-2-propanol mixtures at high alcohol concentration as for the water-1-propanol mixture. Moreover, the structure of this system is slightly different compared with the form of the other potential models, as indicated by the  $g(r)$ . On the other hand, it seems that more alcohols (oxygens) are closer to water molecules, as indicated by the first peak in the  $g(r)$  plot.

Finally, the determination of the temperature-dependent (mass) density is reported for 2-propanol/water mixtures (for 1-propanol/water mixtures, this has already been published<sup>[17]</sup>). As it is evident from Figure 9, density values that belong to the OPLS/AA<sup>[23]</sup> potential are the most distinct – and, as comparison with experimental data for pure 2-propanol shows (Figure 9a), OPLS/AA densities match measured data the least, with about 10% margin. The same has been found for the density and dielectric constant of 1-propanol and its aqueous mixtures<sup>[17]</sup>: the OPLS/AA potential provided worse agreement with experimental densities and dielectric constants than the UA-type force fields. This is, unfortunately, a behavior exactly opposite to what has been established for structural properties in the large part of the present study.



**Figure 8.** Oxygen–oxygen partial radial distribution functions  $g_{O-O}(r)$  (O: alcohol oxygen),  $g_{O-OW}(r)$  (OW: water oxygen), and  $g_{OW-OW}(r)$  for aqueous mixtures with 1-propanol (left) and 2-propanol (right) at 10% molar concentration of the alcohol.



**Figure 9.** Densities of 2-propanol/water mixtures for (a) 100%, (b) 80%, (c) 50%, and (d) 10% alcohol molar concentrations. Experimental data were taken.<sup>[14]</sup> Note the odd behavior observed for the mixture with 50% 2-propanol, which also indicates that most problems occur when the composition is the furthest away from being a neat liquid.

## 5. Summary and Conclusions

New synchrotron X-ray diffraction experiments have been conducted for pure 1- and 2-propanol, as well for their four mixtures with water. Molecular dynamics simulations with both UA and AA type interatomic potentials for the alcohols, and using the TIP4P/ε water potential model, have been performed. X-ray and neutron-weighted total structure factors were determined from the simulation results and compared to diffraction data in the reciprocal space. Densities were also calculated from the molecular dynamics trajectories and compared to available experimental data.

An optimization method<sup>[24]</sup> for obtaining an improved set of potential parameters was applied to the pure 2-propanol OPLS/UA force field.<sup>[22]</sup> The new parameters were validated against the experimental dielectric constant, density and surface tension. Although concerning most of these properties the modified potential does represent an improvement, the calculated total structure factors are not (or, at most, hardly at all) different from those obtained from the original parametrization.<sup>[22]</sup>

It has been shown that the X-ray weighted structure factors may be represented by the UA force field (both the original and the new UA) satisfactorily for the pure alcohols; agreement for the aqueous mixtures may be termed “qualitative” in the best cases. That is, for calculating structural properties, the UA-type potentials are not applicable for 1-propanol/water and 2-propanol/water mixtures. On the other hand, the OPLS/AA force field<sup>[23]</sup> has proven to be consistently better in comparison with the X-ray diffraction data; also, for the reproduction of neutron diffraction data, only the all atom type potential is applicable. Therefore concerning structural properties, the application of the OPLS/AA alcohol potential, in combination with the TIP4P/ε water model,<sup>[24]</sup> is recommended.

The AA- and UA-type potentials produce similar O–O partial radial distribution functions for 1-propanol and its mixtures with water. However, for 2-propanol and its mixtures, only PRDF-s originating to the OPLS/AA potential model may be trusted.

Finally, it has to be admitted that significant differences in terms of the capabilities of the various force fields has been shown for the dielectric constant<sup>[17]</sup> and density (Figure 9): in both cases, the modified UA force field has proven to be better than either the AA,<sup>[23]</sup> or the original UA<sup>[22]</sup> models. That is, even though the structure of 1-propanol/water and 2-propanol/water liquid mixtures is well described by the all atom type potentials applied here, the OPLS/AA force field<sup>[23]</sup> still is not capable of capturing the overall behavior of these systems. In other words, depending on the specific target property, one should, unfortunately, pick different potential models for molecular simulations propanol/water mixtures.

## Acknowledgments

László Temleitner and László Pusztai were supported by the National Research, Development and Innovation Office (Hungary), via grant No. SNN116198. Hector Dominguez and José G. Méndez-Bermúdez acknowledge supercomputer facilities of DGTIC-UNAM (México) grant LANCAD-UNAM-DGTIC-238. László Temleitner is grateful to the János Bolyai Research Scholarship of the Hungarian Academy of Sciences. Synchrotron radiation experiments were performed with the approval of the Japan Synchrotron Radiation Research Institute (JASRI) (Proposal No. 2013A1083). José G. Méndez-Bermúdez acknowledges Programa para el Desarrollo Profesional (PRODEP) for financial support. Hector Domínguez acknowledges DGAPA-UNAM grant IN102017.

## Conflict of Interest

The authors declare no conflict of interest

## Keywords

force field parametrization, liquid structure, molecular dynamics, X-ray diffraction

Received: May 14, 2018  
Revised: August 30, 2018  
Published online: October 7, 2018

- [1] A. Mikusińska-Planner, *Acta Crystallogr. A: Cryst. Phys. Diffr. Theor. Gen. Crystallogr.* **1977**, *33*, 433.
- [2] K. S. Vahvaselkä, R. Serimaa, M. Torkkeli, *J. Appl. Crystallogr.* **1995**, *28*, 189.
- [3] I. Akiyama, M. Ogawa, K. Takase, T. Takamuku, T. Yamaguchi, N. Ohtori, *J. Solution Chem.* **2004**, *33*, 797.
- [4] M. Tomšič, A. Jamnik, G. Fritz-Popovski, O. Glatzer, L. Vlček, *J. Phys. Chem. B* **2007**, *111*, 1738.

- [5] A. Sahoo, S. Sarkar, P. S. R. Krishna, V. Bhagat, R. N. Joarder, *Pramana* **2008**, *71*, 133.
- [6] A. Sahoo, S. Sarkar, V. Bhagat, R. N. Joarder, *J. Phys. Chem. A* **2009**, *113*, 5160.
- [7] P. Sillrén, J. Swenson, J. Mattsson, D. Bowron, A. Matic, *J. Chem. Phys.* **2013**, *138*, 214501.
- [8] L. A. Rodríguez Palomino, G. J. Cuello, A. Stunault, J. Dawidowski, *Philos. Mag.* **2015**, *96*, 816.
- [9] P. Zetterström, U. Dahlborg, W. S. Howells, *Mol. Phys.* **1994**, *81*, 1187.
- [10] A. Sahoo, S. Sarkar, P. S. R. Krishna, R. N. Joarder, *Pramana* **2010**, *74*, 765.
- [11] H. Hayashi, K. Nishikawa, T. Iijima, *J. Phys. Chem.* **1990**, *94*, 8334.
- [12] T. Takamuku, H. Maruyama, K. Watanabe, T. Yamaguchi, *J. Solution Chem.* **2004**, *33*, 641.
- [13] T. Takamuku, K. Saisho, S. Nozawa, T. Yamaguchi, *J. Mol. Liq.* **2005**, *119*, 133.
- [14] J. McGregor, R. Li, J. A. Zeitler, C. D'Agostino, J. H. Collins, M. D. Mantle, H. Manyar, J. D. Holbrey, M. Falkowska, T. G. A. Youngs, C. Hardacre, E. H. Stitt, L. F. Gladden, *Phys. Chem. Chem. Phys.* **2015**, *17*, 30481.
- [15] H. J. Tong, J. Y. Yu, Y. H. Zhang, J. P. Reid, *J. Phys. Chem. A* **2010**, *114*, 6795.
- [16] A. Vrhovsek, O. Gereben, A. Jamnik, L. Pusztai, *J. Phys. Chem. B* **2011**, *115*, 13473.
- [17] J. G. Méndez-Bermúdez, H. Dominguez, L. Pusztai, S. Guba, B. Horváth, I. Szalai, *J. Mol. Liq.* **2016**, *219*, 354.
- [18] S. Kohara, K. Suzuya, Y. Kashihara, N. Matsumoto, N. Umesaki, I. Sakai, *Nucl. Instrum. Methods Phys. Res. A* **2001**, *467*, 1030.
- [19] S. Kohara, M. Itou, K. Suzuya, Y. Inamura, Y. Sakurai, Y. Ohishi, M. Takata, *J. Phys.: Condens. Matter* **2007**, *19*, 506101.
- [20] a) <http://www.gromacs.org>; b) D. van der Spoel, E. Lindahl, B. Hess, G. Groenhof, A. E. Mark, H. J. C. Berendsen, *J. Comput. Chem.* **2005**, *26*, 1701.
- [21] F. Fuentes-Azcatl, J. Alejandro, *J. Phys. Chem. B* **2014**, *118*, 1263.
- [22] W. L. Jorgensen, *J. Phys. Chem.* **1986**, *90*, 1276.
- [23] W. L. Jorgensen, D. S. Maxwell, J. Tirado-Rives, *J. Am. Chem. Soc.* **1996**, *118*, 11236.
- [24] F. J. Salas, G. A. Méndez-Maldonado, E. Nuñez-Rojas, G. E. Aquilar-Pineda, H. Dominguez, J. Alejandro, *J. Chem. Theory Comput.* **2015**, *11*, 683.
- [25] H. J. C. Berendsen, J. P. M. Postma, W. F. van Gunsteren, A. DiNola, J. R. Haak, *J. Chem. Phys.* **1984**, *81*, 3684.
- [26] N. V. Sastry, S. R. Patel, *J. Chem. Thermodyn.* **2000**, *32*, 1669.
- [27] G. Vazquez, E. Alvarez, J. M. Navaza, *J. Chem. Eng. Data* **1995**, *40*, 611.
- [28] F. M. Pang, C. E. Seng, T. T. Teng, W. H. Ibrahim, *J. Mol. Liq.* **2007**, *136*, 71.
- [29] D. Waasmaier, A. Kirfel, *Acta Crystallogr.* **1995**, *A51*, 416.
- [30] V. F. Sears, *Neutron News* **1992**, *3*, 26.